

“Staged” Strong Coupling Strategy for Fluid-Structure Interactions: Supersonic Flow Test Problem

Rodrigo R. PAZ - Lisandro D. DALCÍN

Mario A. STORTI - Norberto M. NIGRO

INTEC-CONICET. Santa Fe, Argentina.
mailto: rodrigop@intec.unl.edu.ar

ENIEF 2006, November 7-10, 2006 - Santa Fe, Argentina.



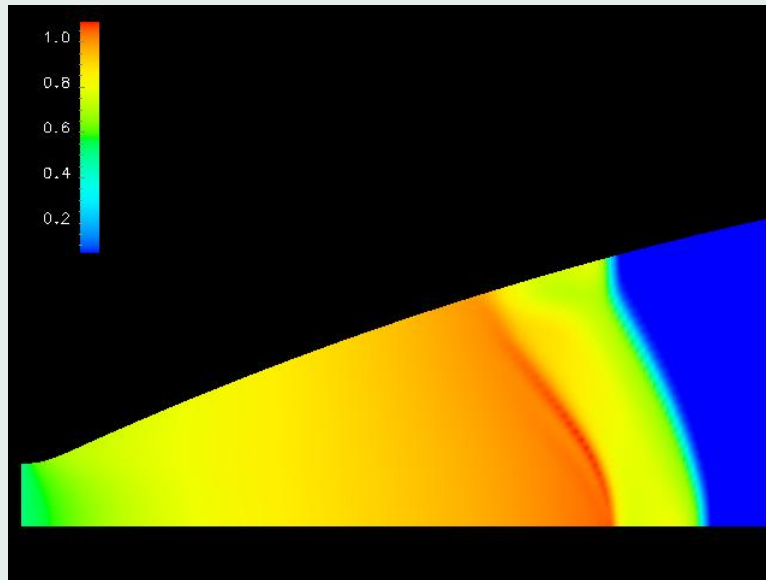
1. Introduction

- One of the three tests proposed by “CIMEC CFD Team” to the **E.S.A. (ESTEC)/OE** to evaluate two CFD codes for its Strong Fluid-Structure Interaction project at supersonic/hypersonic regime:



- ▷ **FINE/HEXA-3D** code by Numeca (C. Hirsch) and OOFELIE by O.E., Belgium.
- ▷ **PETSc-FEM** code from CIMEC-INTEC, Argentina.
(<http://www.cimec.org.ar/petscfem>).

2. Supersonic gas exhaust in the nozzle of the Vulcain rocket engine

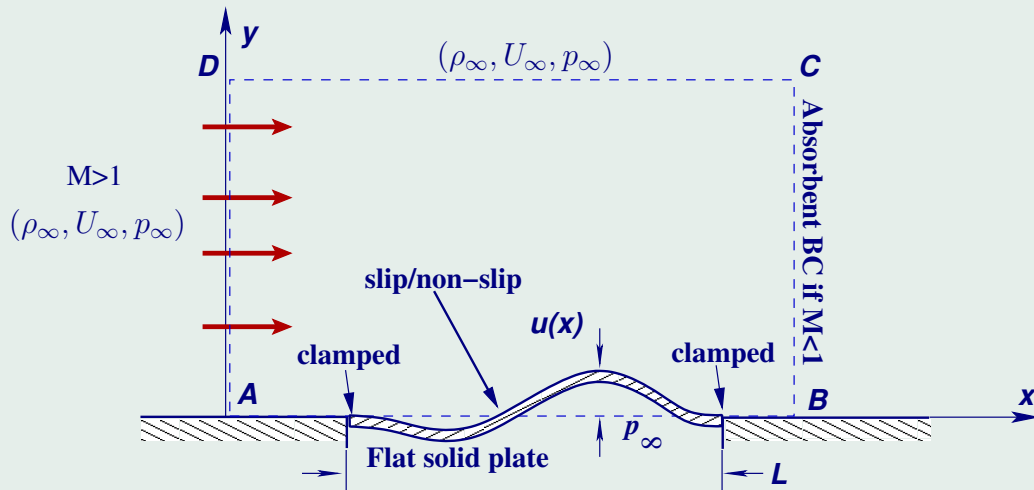


3. Introduction (contd...)

These tests involved different numerical techniques, namely:

- Weak/Strong algorithm to couple Fluid and Structure (separate) codes (Presented in ENIEF 2006 by Mario Storti).
- Parallel code for ‘Inviscid/Viscous Hypersonic Flows’ in ‘Beowulf’ clusters of PC’s.
- The Galerkin/SUPG formulation with added ‘isotropic/anisotropic shock capturing’ operators.
- ALE and ‘Adaptive mesh refinement’ techniques.
- The use of ‘non-reflecting’ boundary conditions on fictitious (and subsonic) walls.
- Finally, a ‘Domain Decomposition’ preconditioner for an efficient solution of the linear system.

4. Test: Flutter of a flat solid plate aligned with the supersonic gas flow



- Viscous/inviscid supersonic flow (EULER/Compressible N-S eqs).
- Thin plate theory for structure: $m\ddot{u}(x, t) + D\frac{\partial^4 u(x, t)}{\partial x^4} = -(p - p_\infty) + f(x, t)$.
- Coupled through **pressure** and **traction** (viscous case only) on interface boundary.
- Undisturbed flow $(\rho, \mathbf{v}, p)_\infty$ is a solution of the problem for zero initial condition (solid problem).

5. Theoretical Approximation: Houbolt's model [1958]

- **Fluid Problem**

$$p - p_{\infty} = C_x \frac{\partial u}{\partial x} + C_t \frac{\partial u}{\partial t},$$
$$C_x = \frac{\rho_{\infty} U_{\infty}^2}{\sqrt{M_{\infty}^2 - 1}}, \quad C_t = \frac{\rho_{\infty} U_{\infty} (M_{\infty}^2 - 2)}{(M_{\infty}^2 - 1)^{3/2}}.$$

- Then, for the **Plate Problem** the deflection becomes

$$m\ddot{u} + D \frac{\partial^4 u}{\partial x^4} = -C_x \frac{\partial u}{\partial x} - C_t \frac{\partial u}{\partial t}.$$

- Using a global basis for displacements

$$u(x) = \sum_{k=1}^N a_k \psi_k(x),$$
$$\psi_k(x) = \frac{4 x(L - x)}{L^2} \sin(k\pi x/L).$$

- The basis functions satisfy the essential boundary conditions for plate equation $u = (\partial u / \partial x) = 0$ at $x = 0, L$.
- Replacing $u(x)$ in the Houbolt's approximation and using Galerkin method

$$M\ddot{a} + Ka + H_x\dot{a} + H_t\ddot{a} = 0,$$

where

$$M_{jk} = \int_0^L m \psi_j(x) \psi_k(x) dx,$$

$$K_{jk} = \int_0^L D \psi_j''(x) \psi_k''(x) dx,$$

$$H_{x,jk} = \int_0^L C_x \psi_j(x) \psi_k'(x) dx,$$

$$H_{t,jk} = \int_0^L C_t \psi_j(x) \psi_k(x) dx.$$

- If the ansatz $\mathbf{a}(t) = \hat{\mathbf{a}}e^{\lambda t}$ is proposed as a solution for the **Plate Problem**, the following eigenvalue problem is stated

$$(\lambda^2 \mathbf{M} + \lambda \mathbf{H}_t + \mathbf{K} + \mathbf{H}_x) \hat{\mathbf{a}} = 0.$$

- Using time and mass non-dimensional parameters

$$N_T = \left(\frac{T_{fl}}{T_{str}} \right)^2 = \frac{L/U_\infty}{\sqrt{mL^4/D}} = \frac{D}{mL^2U_\infty^2} \quad \text{and}$$

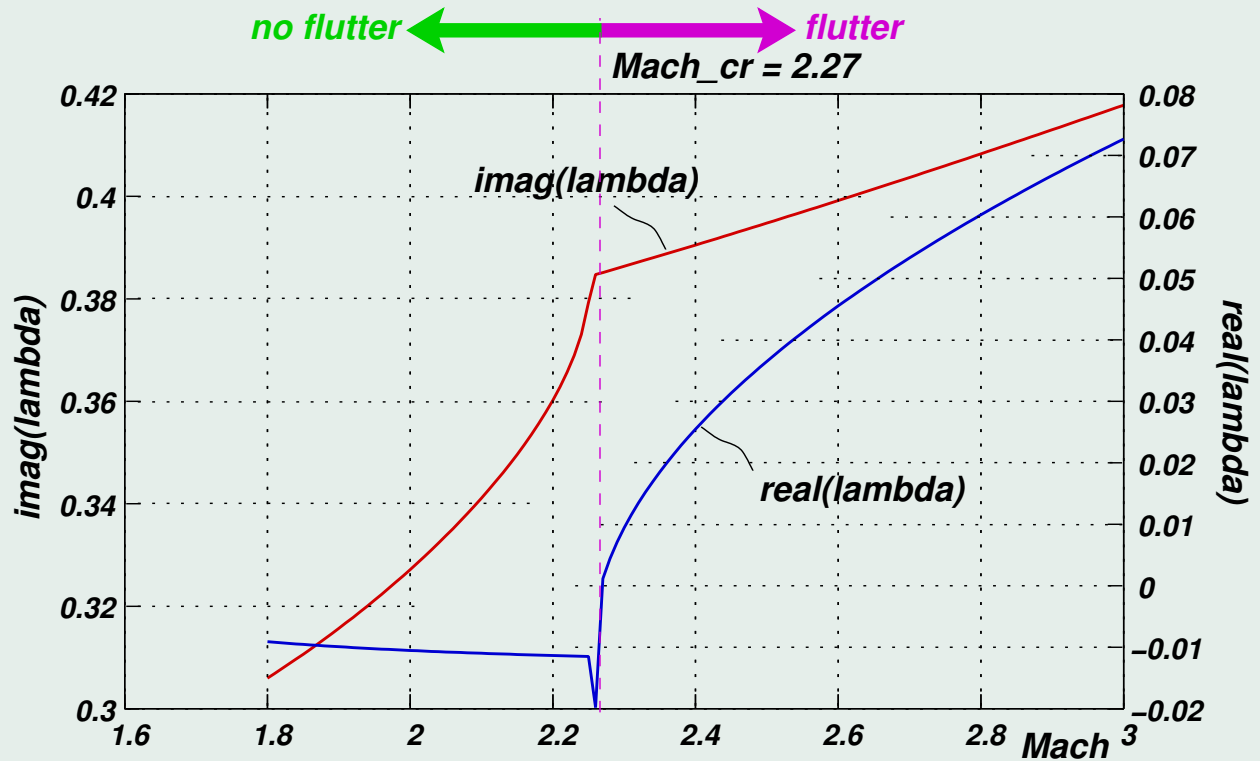
$$N_M = \frac{\rho_\infty L^3}{mL^2} = \frac{\rho_\infty L}{m},$$

the space of parameters is full covered.

- Results for the Houbolt's model

▷ $N = 20$, $N_x = 5000$, and a sweep in M_∞ while keeping constant ρ_∞ , m , L and D ,

i.e., $N_M = \text{cte}$ and $N_T \propto M_\infty^{-2}$.



**Flutter mode appears for $M_{\infty} \geq 2.27$
or $N_T = 4.3438 \times 10^{-5}$ and $N_M = 0.055$.**

6. Staged FSI-FEM Results for flutter

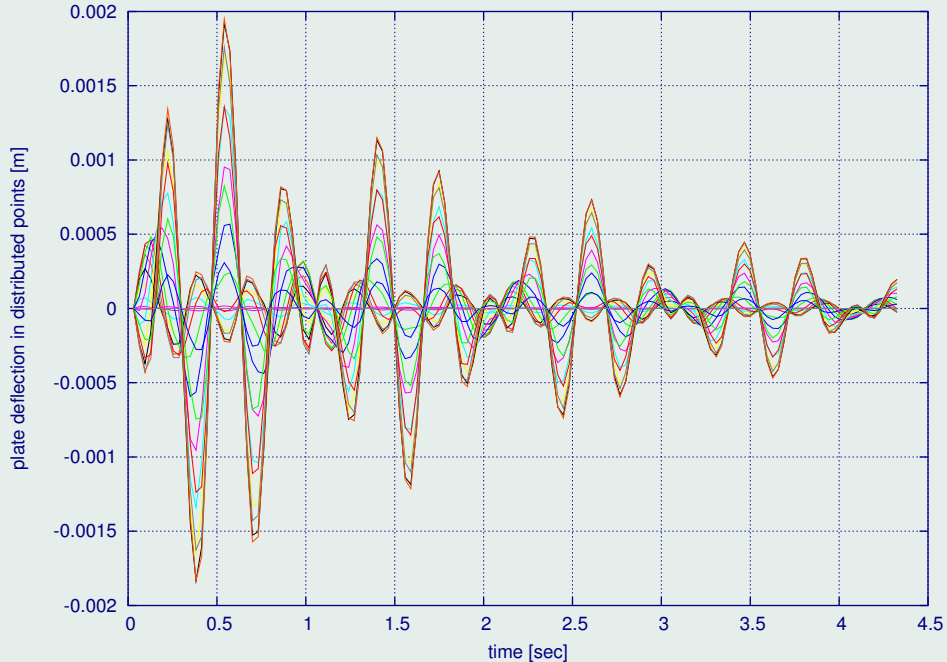


Figure 1: Structure response at $M=1.8$

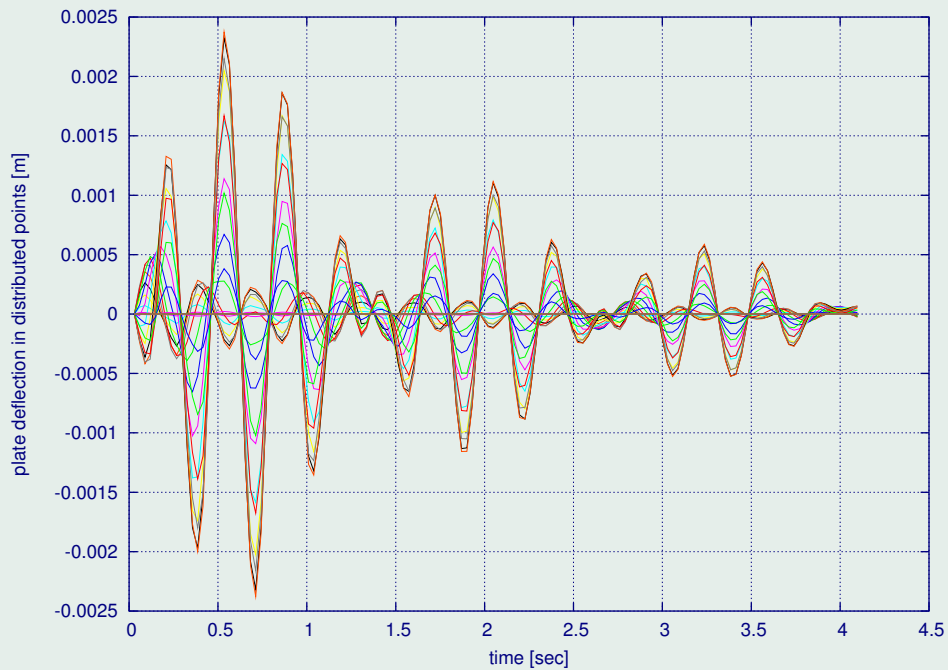


Figure 2: Structure response at $M=2.0$

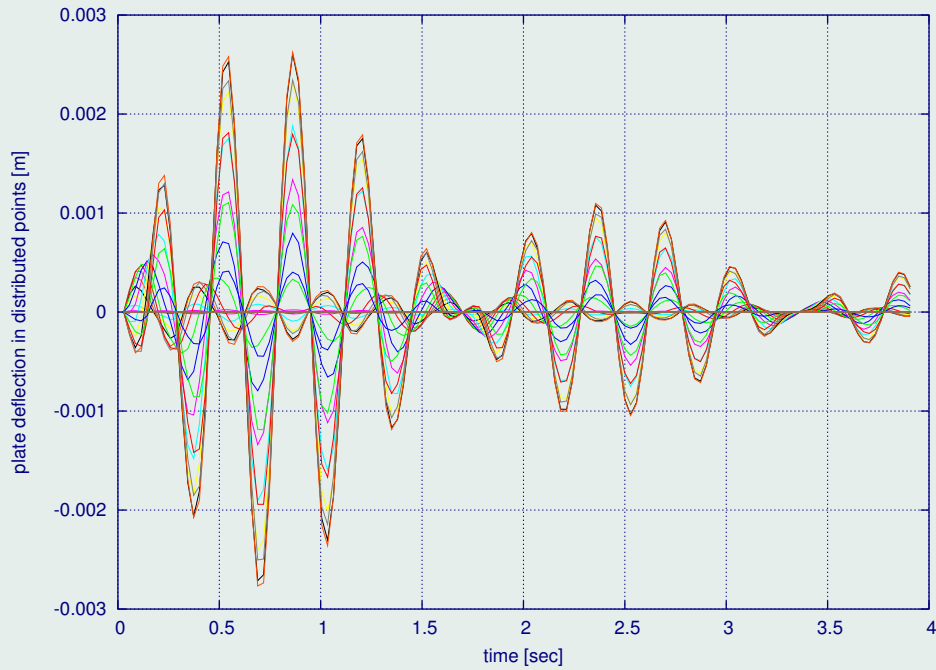


Figure 3: Structure response at $M=2.1$

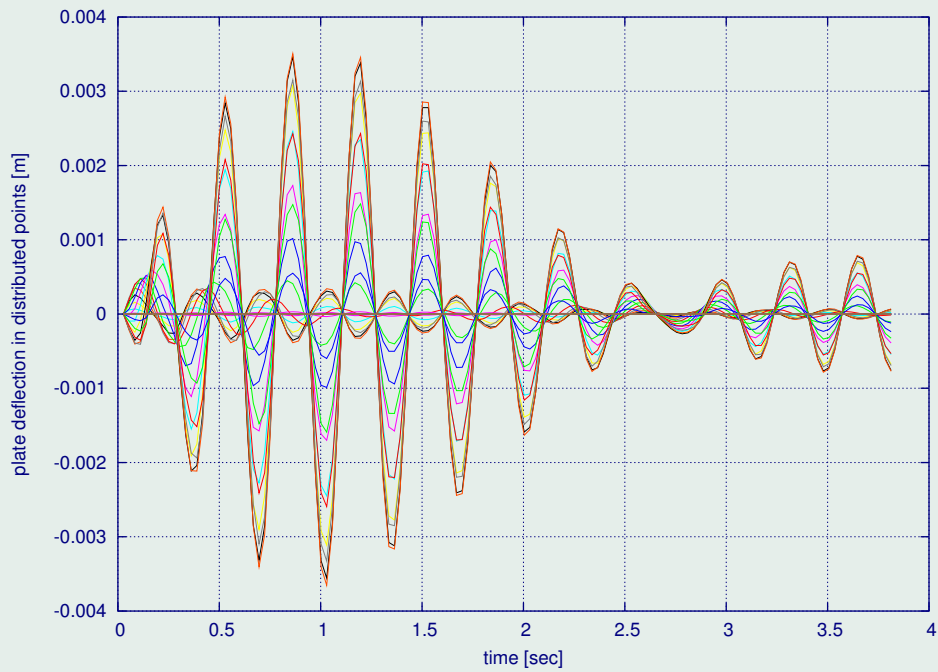


Figure 4: Structure response at $M=2.2$

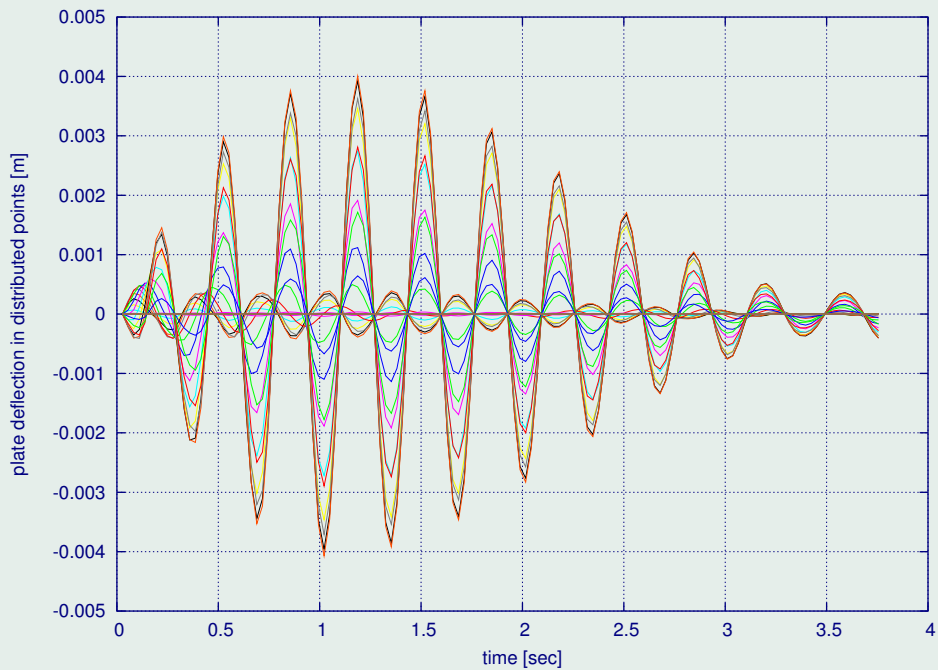


Figure 5: Structure response at $M=2.225$

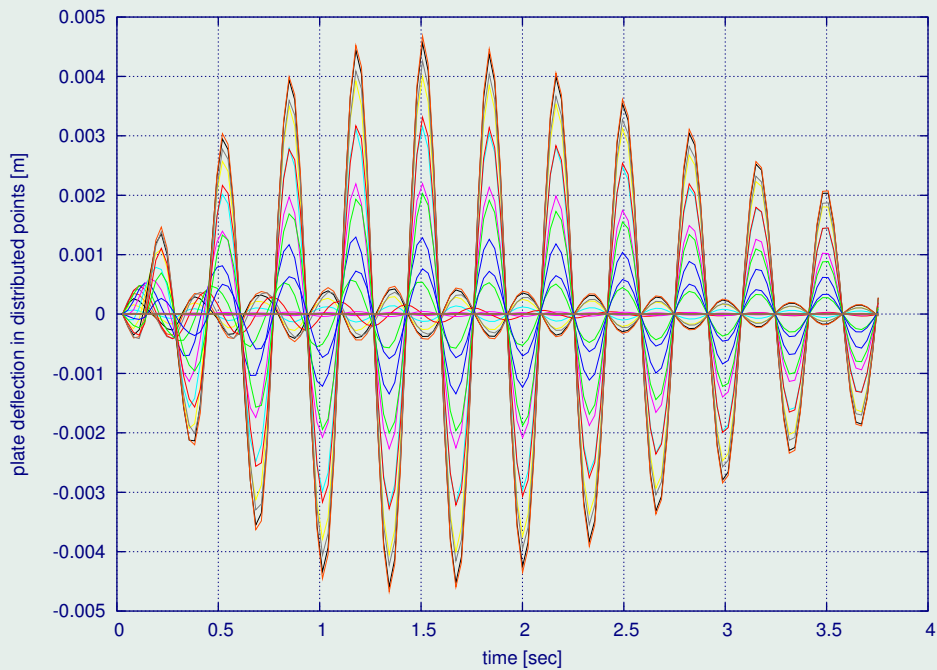


Figure 6: Structure response at $M=2.25$

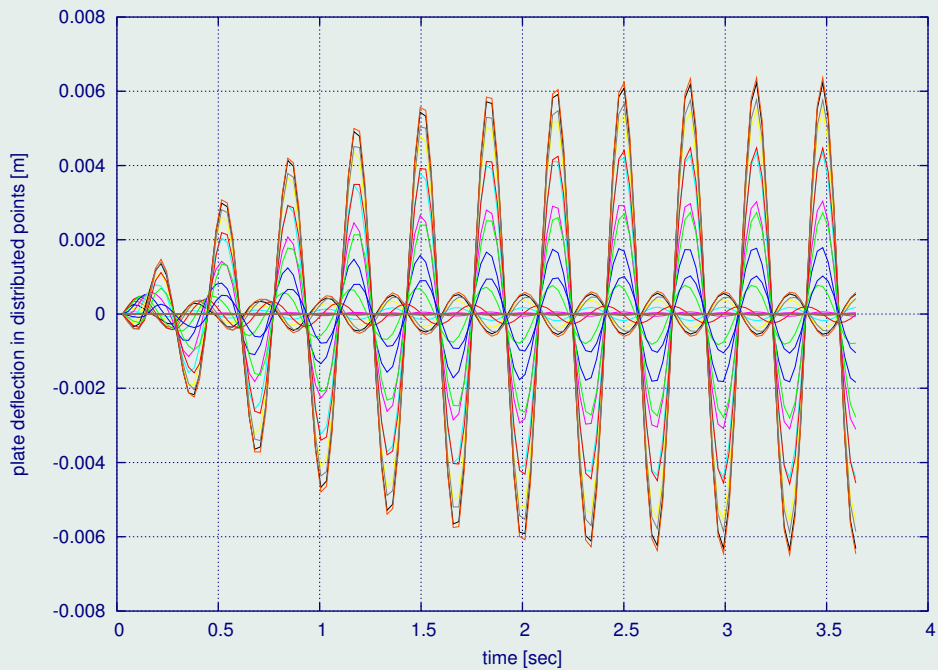


Figure 7: Structure response at $M=2.275$

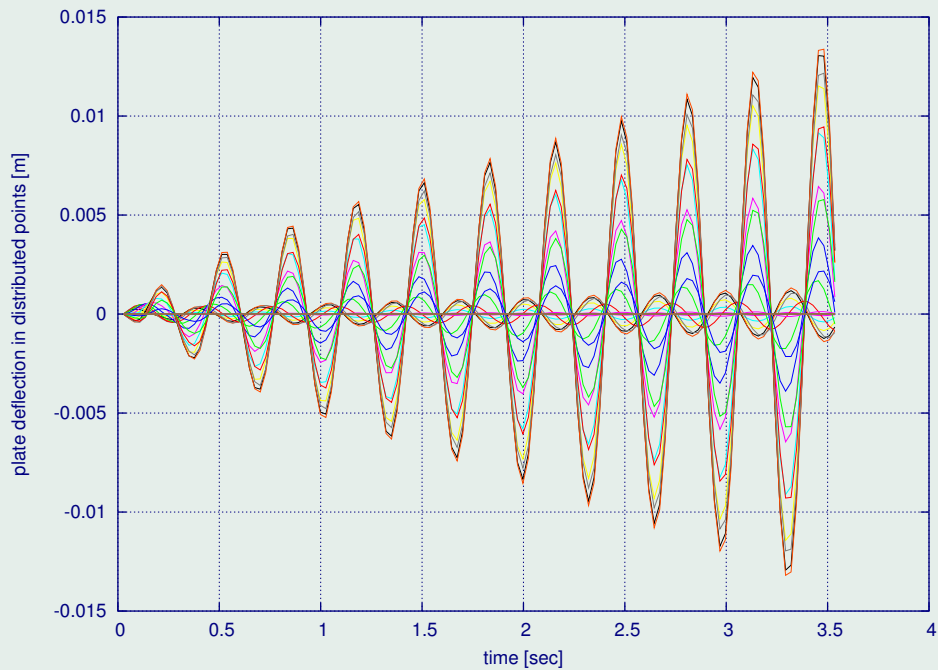


Figure 8: Structure response at $M=2.3$

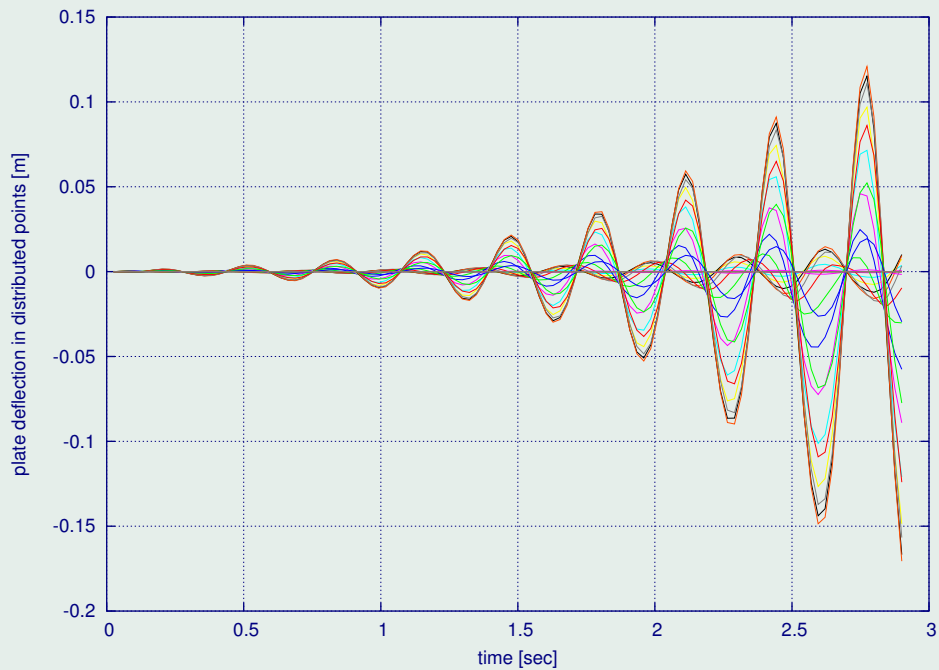


Figure 9: Structure response at $M=2.5$

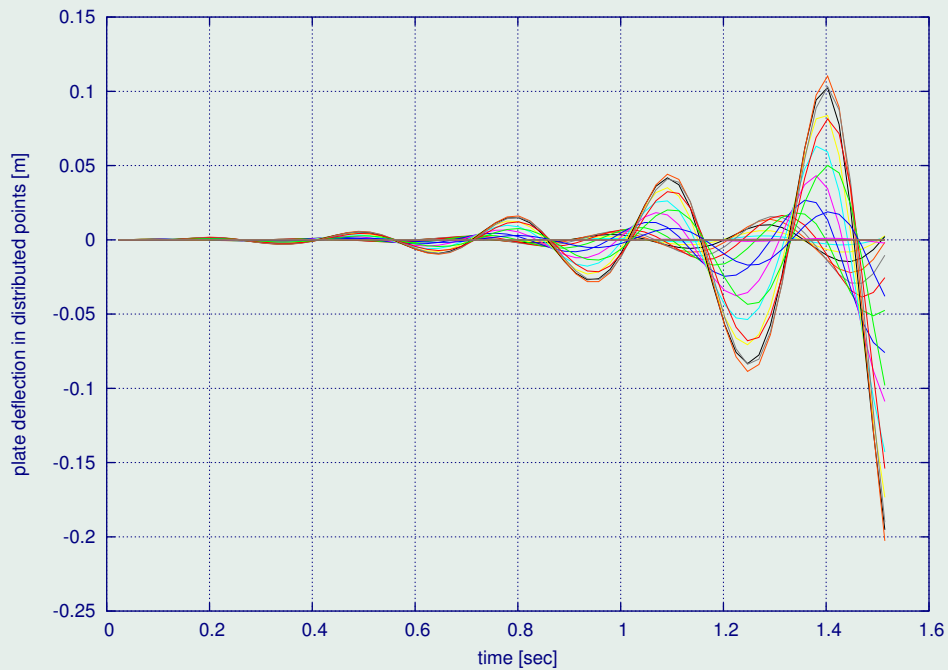


Figure 10: Structure response at $M=3.0$

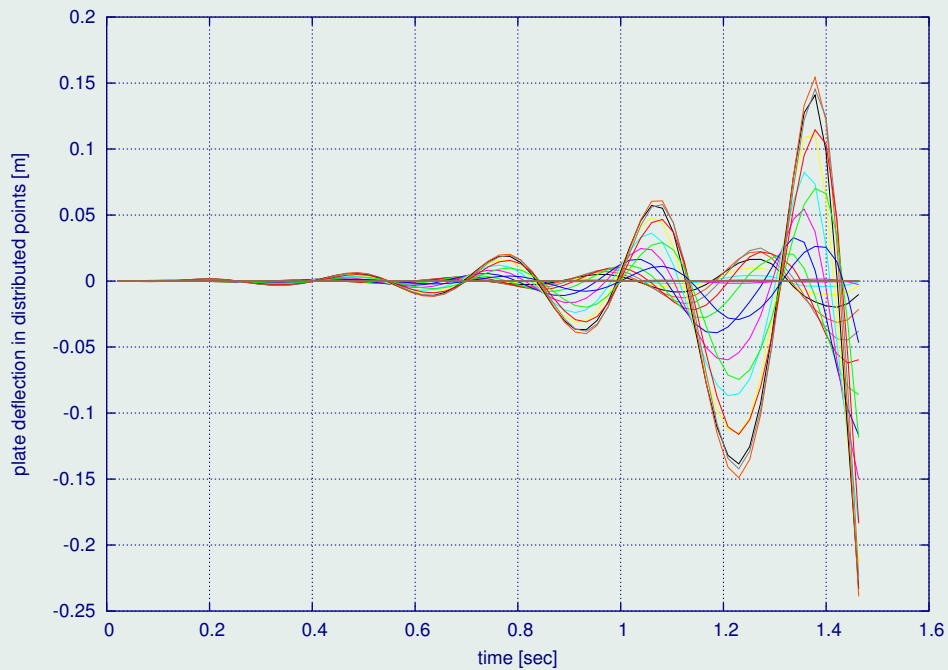


Figure 11: Structure response at $M=3.2$

7. Flat plate in Flutter

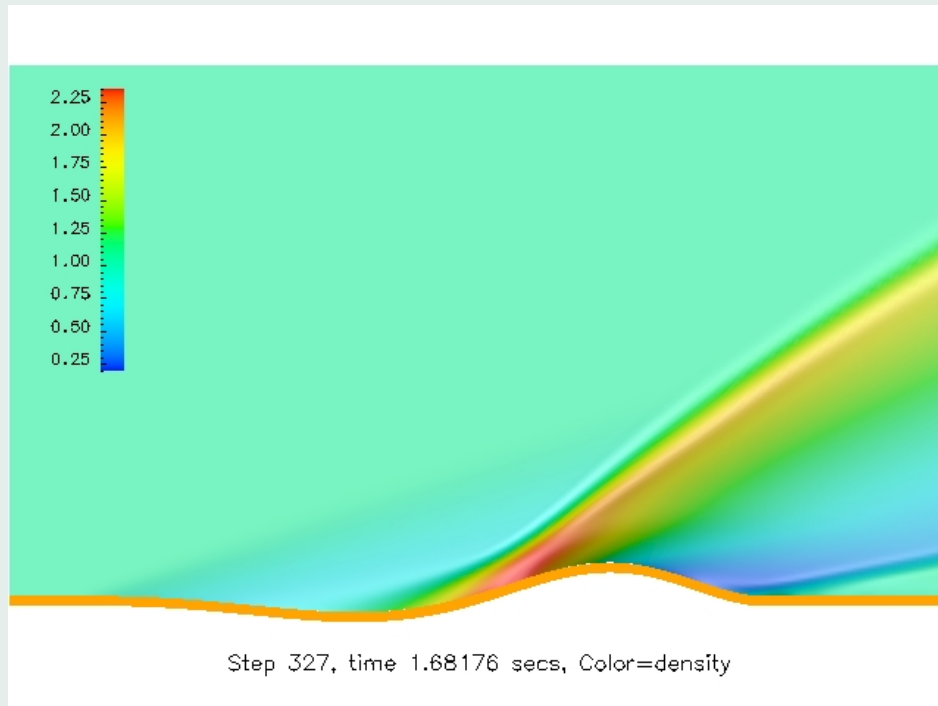


Figure 12: Fluid and Solid fields at $M_{\infty} = 3.2$

- Being

$$\frac{N_M}{N_T M_\infty^2} = \frac{\rho_\infty L^3 c_\infty}{D},$$

sweeps in N_T , N_M and M_∞ estimated the **flutter region** as

$$\frac{N_M}{N_T M_\infty^2} < 200 \text{ no flutter for any Mach number,}$$

$$\frac{N_M}{N_T M_\infty^2} > 300 \text{ flutter for the lowest Mach considered } (M_\infty \geq 1.8).$$

- If $(\partial \mathbf{u} / \partial \mathbf{t})$ is neglected (i.e., characteristic times of struct. are much lower than those of the fluid, $N_T \ll 1$),

$$\det(\bar{\lambda}^2 \bar{M} + K + H_x) = 0$$

$$\bar{\lambda} = \sqrt{m} \lambda,$$

$$\bar{M} = \frac{1}{\sqrt{m}} M,$$

the coefficients in \bar{M} , K , H_x do not depend on m , neither do the eigenvalues of equation, then the λ eigenvalues are of the form

$$\lambda_j = \frac{\bar{\lambda}_j}{\sqrt{m}},$$

with $\bar{\lambda}$ not depending on m . This means that the sign of the real part of the λ is independent of m .

8. Stability of the staged algorithm (flutter region, contd...)

$$U_{\infty} = M_{\infty} = 2$$

$$t = 0.06$$

$$\nu = 0.33$$

$$m = 0.002$$

$$E = 39.6$$

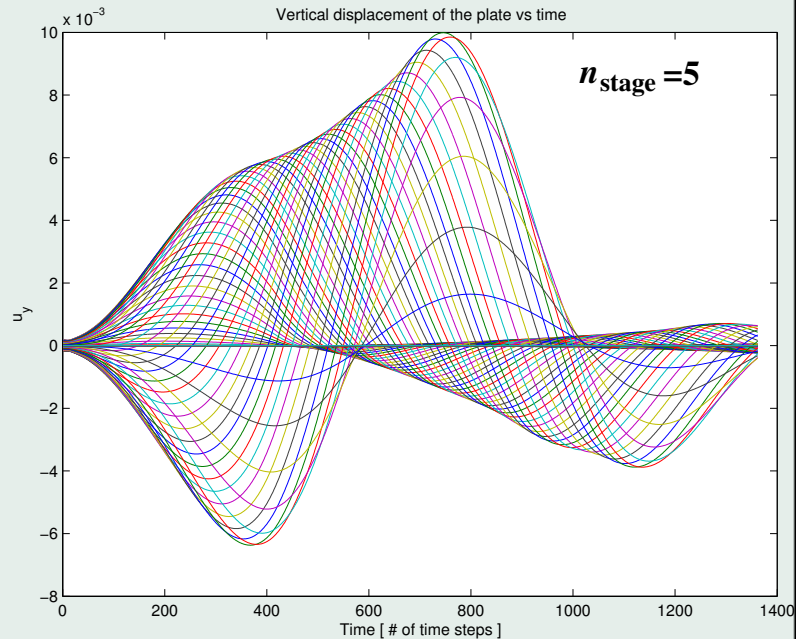
$$D = 8.010^{-4}$$

$$N_T = \frac{D}{mL^2U_{\infty}^2} = 0.025$$

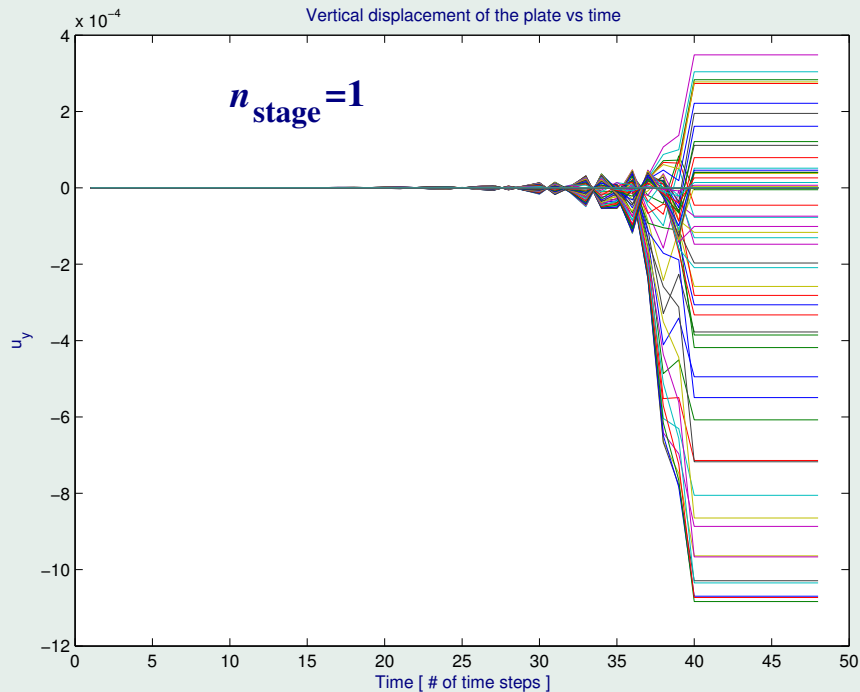
$$N_M = \frac{\rho_{\infty}L}{m} = 1000.0$$

$$\frac{N_M}{N_T M_{\infty}^2} = 10000 > 300$$

(i.e., inside the **flutter region**)



9. Stability of the staged algorithm (flutter region, contd...)



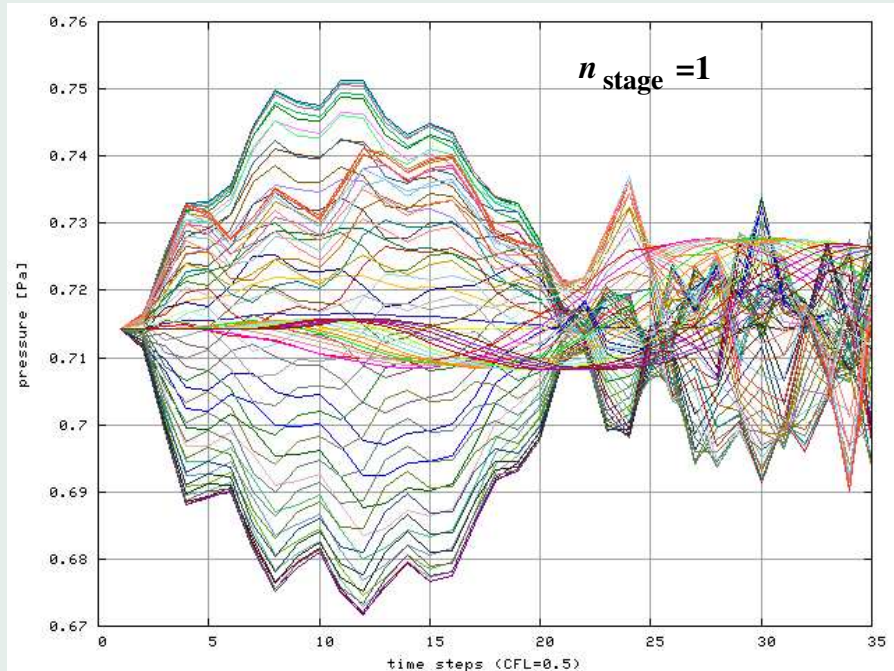
10. Stability of the staged algorithm outside the flutter region, i.e. $N_M/(N_TM_\infty^2) \ll 200$

- $N_M/(N_TM_\infty^2)$ do not depend on plate density m

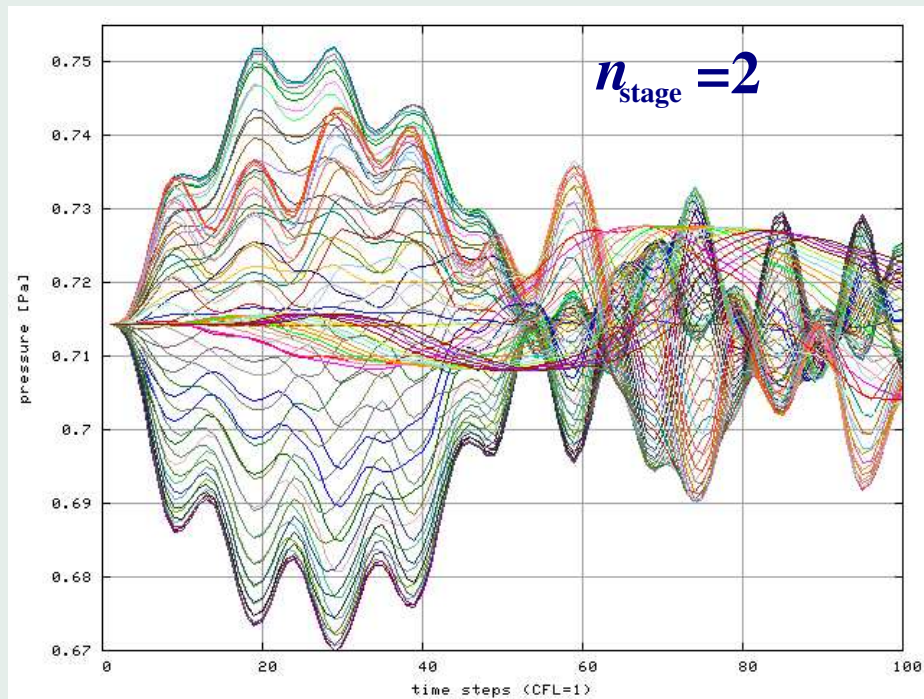
with $m = 0.00135$

$$\frac{N_M}{N_TM_\infty^2} = 12 < 200$$

(i.e., outside the
flutter region)



11. Stability of the staged algorithm outside the flutter region (contd...)



12. Stability of the staged algorithm (flutter region)

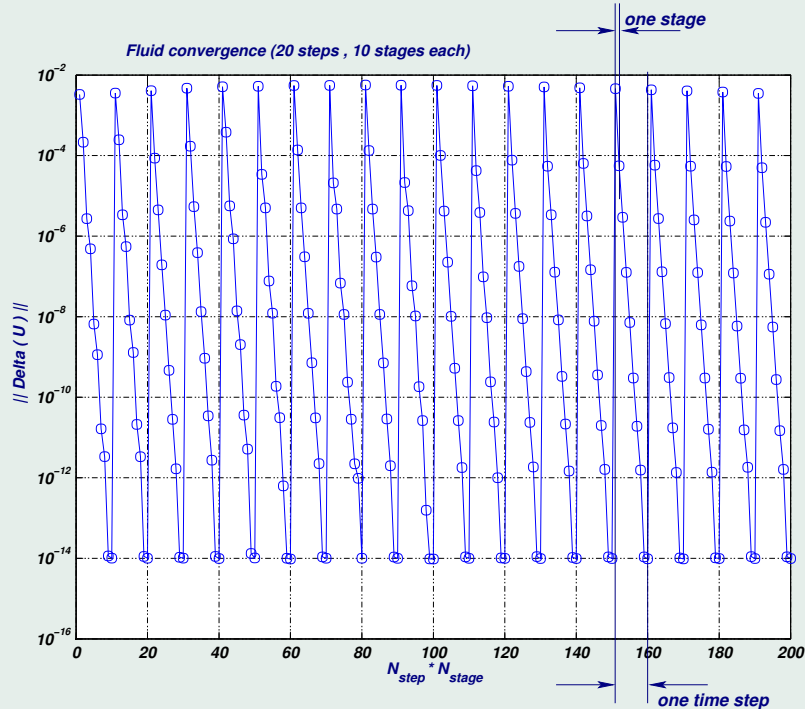


Figure 13: Convergence of fluid state in stage loop

13. Conclusions

- “Staged” strategy provides a smooth blending between weak coupling and strong coupling.
- Moderately coupled problems that can not be treated with the pure weak coupling approach, can be solved with the staged algorithm using few stages per time step.
- The elastic flat plate problem is geometrically simple, but gives physical insight in the flutter phenomena, and was very useful in testing the proposed algorithm in a wide range of non-dimensional parameters.

Acknowledgment

- This work has received financial support from Consejo Nacional de Investigaciones Científicas y Técnicas (CONICET, Argentina), Agencia Nacional de Promoción Científica y Tecnológica (ANPCyT) and Universidad Nacional del Litoral (UNL) through grants PIP 0198/98, PIP 02552/00, PIP 5271/05, PICT Lambda 12-14573/2003, PME 209/2003.
- We made extensive use of freely distributed software such as **GNU/Linux OS, MPI, PETSc, Newmat, Metis, Octave, the CGAL geometrical library, OpenDX** and many others.