

CAPILLARY PRESSURE SENSITIVITY IN MODELING CO₂ INJECTION

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Abstract. Carbon dioxide (CO₂) sequestration consists in injecting and storage the gas into a geologic formation as a means of mitigating the greenhouse effect. Among the storage sites, saline aquifers are very promising because of their large capacity and wide availability. In this work, we perform a sensitivity analysis on the response of a flow simulator to study the impact of capillary pressure in CO₂ injection and storage. Capillary pressure is represented as a exponential function of CO₂ saturation. The simultaneous flow of CO₂ and brine in porous media is based on the well known Black-Oil formulation, applied to two-phase fluid flow. It considers that CO₂ may dissolve in the brine but the brine is not allowed to vaporize into the CO₂ phase. This formulation uses the PVT data as a simplified thermodynamic model. The numerical solution is obtained applying the public-domain software BOAST, which solves the differential equations by finite differences with the IMPES algorithm. Besides, we build a suitable geological model based on fractal porosity and clay content and taking into account the variation of properties with pore pressure and CO₂ saturation. This model also includes embedded mudstone layers of very low permeability that accumulate CO₂ but also allow its upward migration. Relative permeability and capillary pressure as functions of CO₂ saturation are taken from published data, using the Leverett *J*-function to scale capillary pressure values. The numerical results show that the fluid simulator is able to represent the CO₂ injection and storage. As injection proceeds, part of the injected fluid migrates upwards across the mudstone layers. The simulations are performed using different capillary pressure curves. When capillary pressure is higher, CO₂ upward migration is slower and, consequently, different zones of accumulations below the layers and high CO₂ saturations levels between layers are obtained. Thus, a precise estimation of the capillary pressure function is needed to perform realistic long term predictions of the spatial-time CO₂ distribution in the underground.

1 INTRODUCTION

Numerical modeling of CO₂ injection and storage is an important tool to analyze the effectiveness of this procedure in reducing the amount of greenhouse gases in the atmosphere (Arts et al., 2008). In fact, there is no experience in the long term behavior of stored CO₂. The first industrial CO₂ injection project begins in 1996 in the Sleipner gas field in Norway. In this project, CO₂ separated from natural gas is being injected in the Utsira formation, a high permeable sandstone, with several mudstone layers that limit the vertical motion of the CO₂ (Arts et al., 2008)-(Chadwick et al., 2005).

We introduce a methodology to simulate the CO₂ injection and storage in saline aquifers, combined with a petrophysical model of the Utsira formation. The petrophysical model is based on initial porosity and clay content and considers the variation of properties with pore pressure and saturation (Carcione et al., 2003). Therefore, these properties are time dependent due to the CO₂ injection, but they change at a much slower rate than pressure and saturations. As a consequence, we have two time scales, using a much larger time step to update petrophysical properties than to run the flow simulator. The initial porosity has a fractal spatial distribution (Frankel and Clayton, 1986) and the permeability is assumed to be anisotropic and is obtained from first principles as a function of porosity and grain sizes. This geological model is able to simulate embedded mudstone layers of very low permeability.

The Black-Oil formulation for two-phase flow in porous media (Aziz and Settari, 1985) is applied to represent the CO₂ - brine flow. PVT data are computed from equations of state using Hassanzadeh correlations (Hassanzadeh et al., 2008). The numerical solution is obtained with the public-domain software BOAST (Fanchi, 1997), which solves the differential equations by finite differences with the IMPES algorithm (Implicit Pressure Explicit Saturation) (Aziz and Settari, 1985). The basic idea of IMPES is to obtain a single pressure equation combining the flow equations. Once pressure is implicitly computed for the new time, saturation is updated explicitly.

Multiphase flow functions, Brine and CO₂ relative permeabilities and capillary pressure, are taken from published experimental data (Bachu and Bennion, 2008). The relative permeabilities are adjusted with power-law models (Savioli and Bidner, 2005), and capillary pressure is adapted to the Utsira formation applying the Leverett J -function (Leverett, 1941). Besides, we test different capillary pressure curves to perform a sensitivity analysis of the simulator results.

2 CO₂ - BRINE FLOW MODEL IN POROUS MEDIA

The simultaneous flow of brine and CO₂ in porous media is described by the well-known Black-Oil formulation (used in petroleum simulation) applied to two-phase, two component fluid flow (Aziz and Settari, 1985). In this approach, brine is identified with oil and CO₂ with gas, therefore, CO₂ component may dissolve in the brine phase but the brine is not allowed to vaporize into the CO₂ phase.

The Black-Oil formulation uses, as a simplified thermodynamic model, the PVT data defined as

- R_s : CO₂ solubility in brine
- B_g : CO₂ formation volume factor
- B_b : brine formation volume factor

The conversion of compositional data from equations of state into the Black-Oil PVT data is performed applying an algorithm developed by [Hassanzadeh et al. \(2008\)](#),

$$\begin{aligned} \bullet R_s &= \frac{\tilde{\rho}_b^{SC} \chi_g}{\tilde{\rho}_g^{SC} (1 - \chi_g)} \\ \bullet B_b &= \frac{\rho_b^{SC}}{\rho_b (1 - \omega_g)}, \end{aligned}$$

where, ρ_b^{SC} and ρ_g^{SC} are the brine and CO₂ molar densities at standard conditions and χ_g and ω_g are the CO₂ mole and mass fraction in the brine phase.

The differential equations, obtained by combining the mass conservation equations with Darcy's empirical law, are

$$\begin{aligned} \nabla \cdot \left(\underline{\kappa} \left(\frac{k_{rg}}{B_g \eta_g} (\nabla p_g - \rho_g g \nabla D) + \frac{R_s k_{rb}}{B_b \eta_b} (\nabla p_b - \rho_b g \nabla D) \right) \right) + \frac{q_g}{\rho_g^{SC}} \\ = \frac{\partial \left[\phi \left(\frac{S_g}{B_g} + \frac{R_s S_b}{B_b} \right) \right]}{\partial t}, \end{aligned} \quad (1)$$

$$\nabla \cdot \left(\underline{\kappa} \frac{k_{rb}}{B_b \eta_b} (\nabla p_b - \rho_b g \nabla D) \right) + \frac{q_b}{\rho_b^{SC}} = \frac{\partial \left[\phi \frac{S_b}{B_b} \right]}{\partial t}, \quad (2)$$

where g , b denote gas and brine phases, respectively. The unknowns for the Black-Oil model are the fluid pressures p_β and saturations S_β ($\beta = b, g$). Also ρ_β is density, q_β injection mass rate per unit volume, $k_{r\beta}$ relative permeability and η_β viscosity. Finally ϕ is porosity and $\underline{\kappa}$ is the absolute permeability tensor.

Two algebraic equations relating the saturations and pressures complete the system:

$$S_b + S_g = 1, \quad p_g - p_b = P_C(S_b), \quad (3)$$

where P_C is the capillary pressure.

The numerical solution was obtained employing the public domain software BOAST ([Fanchi, 1997](#)). BOAST solves the differential equations using IMPES (IMplicit Pressure Explicit Saturation), a finite difference technique ([Aziz and Settari, 1985](#)). Finite differences is the standard in commercial reservoir simulators, and the improved versions use both structured and unstructured grids with local refinements to accurately represent reservoir geometry. The basic idea of IMPES is to obtain a single pressure equation by a combination of the flow equations. Once pressure is implicitly computed for the new time, saturation is updated explicitly. We briefly describe IMPES for these particular system (1), (2), (3). The first step is to obtain the pressure equation, combining flow equations: Eq. (1) multiplied by B_g and Eq.(2) multiplied by $(B_b - R_s B_g)$ are added. In this way, the right side of the resulting equation is:

$$B_g \frac{\partial \left[\phi \left(\frac{S_g}{B_g} + \frac{R_s S_b}{B_b} \right) \right]}{\partial t} + (B_b - R_s B_g) \frac{\partial \left[\phi \frac{S_b}{B_b} \right]}{\partial t}$$

Using the chain rule to expand the time derivatives and after some computations and rearrangements, the right side becomes:

$$\phi \left[\frac{1}{\phi} \frac{d\phi}{dp_b} + S_g \left(-\frac{1}{B_g} \frac{dB_g}{dp_b} \right) + S_b \left(-\frac{1}{B_b} \frac{dB_b}{dp_b} + \frac{B_g}{B_b} \frac{dR_s}{dp_b} \right) \right] \frac{\partial p_b}{\partial t},$$

where all time derivatives of saturation disappear.

Defining the following approximate compressibilities,

- Formation compressibility: $c_f = \frac{1}{\phi} \frac{d\phi}{dp_b}$
- CO₂ compressibility: $c_g = -\frac{1}{B_g} \frac{dB_g}{dp_b}$,
- Brine compressibility: $c_b = -\frac{1}{B_b} \frac{dB_b}{dp_b} + \frac{B_g}{B_b} \frac{dR_s}{dp_b}$,
- Total compressibility: $c_t = c_f + S_g c_g + S_b c_b$,

the right side is simply expressed as,

$$\phi c_t \frac{\partial p_b}{\partial t}.$$

Finally, replacing p_g by $p_b + P_C(S_b)$ in the left side of the combined equation, the following pressure equation in p_b is obtained,

$$\begin{aligned} & B_g \left[\nabla \cdot \left(\frac{\kappa}{B_g \eta_g} \left(\frac{k_{rg}}{B_g \eta_g} (\nabla p_b - \rho_g g \nabla D) + \frac{R_s k_{rb}}{B_b \eta_b} (\nabla p_b - \rho_b g \nabla D) + \frac{k_{rg}}{B_g \eta_g} \nabla P_C \right) \right) \right] \\ & + (B_b - R_s B_g) \left[\nabla \cdot \left(\frac{\kappa}{B_b \eta_b} \left(\frac{k_{rb}}{B_b \eta_b} (\nabla p_b - \rho_b g \nabla D) \right) \right) \right] \\ & + B_g \frac{q_g}{\rho_g^{SC}} + (B_b - R_s B_g) \frac{q_b}{\rho_b^{SC}} = \phi c_t \frac{\partial p_b}{\partial t}. \end{aligned} \quad (4)$$

In BOAST simulator, system (2), (4) is discretized using a block centered grid. The system is linearized evaluating the pressure and saturation dependent functions (PVT parameters, viscosities, relative permeabilities and capillary pressure) in the pressure and saturation values of the previous time step. The pressure equation is solved implicitly, applying a Block Successive Over Relaxation method (BSOR) to compute the linear system solution. The saturation equation is solved explicitly, therefore stability restrictions are considered to select the time step (Savioli and Bidner, 2005).

3 PETROPHYSICAL MODEL OF A SHALY SANDSTONE

The pressure dependence of properties is based on the following relationship between porosity and pore pressure $p(t) = S_b p_b(t) + S_g p_g(t)$,

$$\frac{(1 - \phi_c)}{K_s} (p(t) - p_H) = \phi_0 - \phi(t) + \phi_c \ln \frac{\phi(t)}{\phi_0}, \quad (5)$$

where ϕ_c is a critical porosity, ϕ_0 is the initial porosity at hydrostatic pore pressure p_H and K_s is the bulk modulus of the solid grains (Carcione et al., 2003). The rock is formed with quartz (bulk modulus of 40 GPa) and clay (bulk modulus of 15 GPa). K_s is computed as the arithmetic average of the Hashin Shtrikman upper and lower bounds. The relationship among horizontal permeability, porosity and clay content C is (Carcione et al., 2003),

$$\frac{1}{\kappa_x(t)} = \frac{45(1 - \phi(t))^2}{\phi(t)^3} \left(\frac{(1 - C)^2}{R_q^2} + \frac{C^2}{R_c^2} \right), \quad (6)$$

where R_q and R_c are the average radii of the sand and clay grains. Also, as permeability is anisotropic, we assume the following relationship between horizontal and vertical permeability κ_z (Carcione et al., 2003)

$$\frac{\kappa_x(t)}{\kappa_z(t)} = \frac{1 - (1 - 0.3a) \sin \pi S_b}{a(1 - 0.5 \sin \pi S_b)}, \quad (7)$$

where a is the permeability-anisotropy parameter.

4 MULTIPHASE FLOW FUNCTIONS

Brine and CO₂ relative permeabilities and capillary pressure are functions of CO₂ saturation. We use the experimental data presented in (Bachu and Bennion, 2008). The relative permeability curves are adjusted using the following power-law models, with two parameters, the maximum value, κ_r^* , and the curvature n ((Savioli and Bidner, 2005),

$$\kappa_{rg}(S_g) = \kappa_{rg}^* \left(\frac{S_g - S_{gc}}{1 - S_{gc} - S_{bc}} \right)^{n_g} \quad (8)$$

$$\kappa_{rb}(S_g) = \kappa_{rb}^* \left(\frac{1 - S_g - S_{bc}}{1 - S_{gc} - S_{bc}} \right)^{n_b} \quad (9)$$

where S_{gc} is the CO₂ critical saturation and S_{bc} is the connate brine saturation.

We represent the relationship between capillary pressure and CO₂ saturation by the following exponential function,

$$P_c(S_g) = P_c^* e^{(n_c S_g)}, \quad (10)$$

where P_c^* is the entry (or threshold) capillary pressure, which means the pressure at which the gas phase is sufficiently connected to allow flow.

The data presented in (Bachu and Bennion, 2008) correspond to a sandstone with different petrophysical properties than the Utsira formation. Then we apply the Leverett J -function (Leverett, 1941)

$$J(S_g) = \frac{P_c(S_g)}{\sigma \cos \theta} \sqrt{\frac{\kappa}{\phi}} \quad (11)$$

to scale the capillary pressure curve, using porosity and permeability data of both formations. The same procedure is applied for the mudstone layers.

5 NUMERICAL EXAMPLES

5.1 Aquifer model

We consider a model of the Utsira formation having 1.2 km in the x -direction, 10 km in the y -direction and 0.4 km in the z -direction (top at 0.77 km and bottom at 1.17 km b.s.l.).

The simulation is performed using a mesh with equally-spaced blocks in each direction: $n_x = 300$ in the x -direction, $n_y = 5$ in the y -direction and $n_z = 400$ in the z -direction. Actually the model is 2.5D since the properties are uniform along the y -direction, which has an extension of 10 km.

The pressure-temperature conditions are $T = 31.7z + 3.4$, where T is the temperature (in °C) and z is the depth (in km b.s.l.); $p_H = \rho_b gz$ is the hydrostatic pressure, with $\rho_b = 1040$ kg/m³ the density of brine and g the gravity constant.

The initial porosity ϕ_0 (at hydrostatic pore pressure) for the Utsira sandstone is assumed to have a fractal spatial distribution, based on the so-called von Karman self-similar correlation functions (Frankel and Clayton, 1986). Horizontal and vertical permeabilities are determined by using equations (6) and (7), considering an anisotropy parameter $a = 0.1$ and a fixed clay content $C = 6$ %. The spatial vertical permeability distribution at initial conditions is shown in Figure 1. It can be observed that, within the formation, there are several mudstone layers (low permeability and porosity) which act as barriers to the vertical motion of the CO₂.

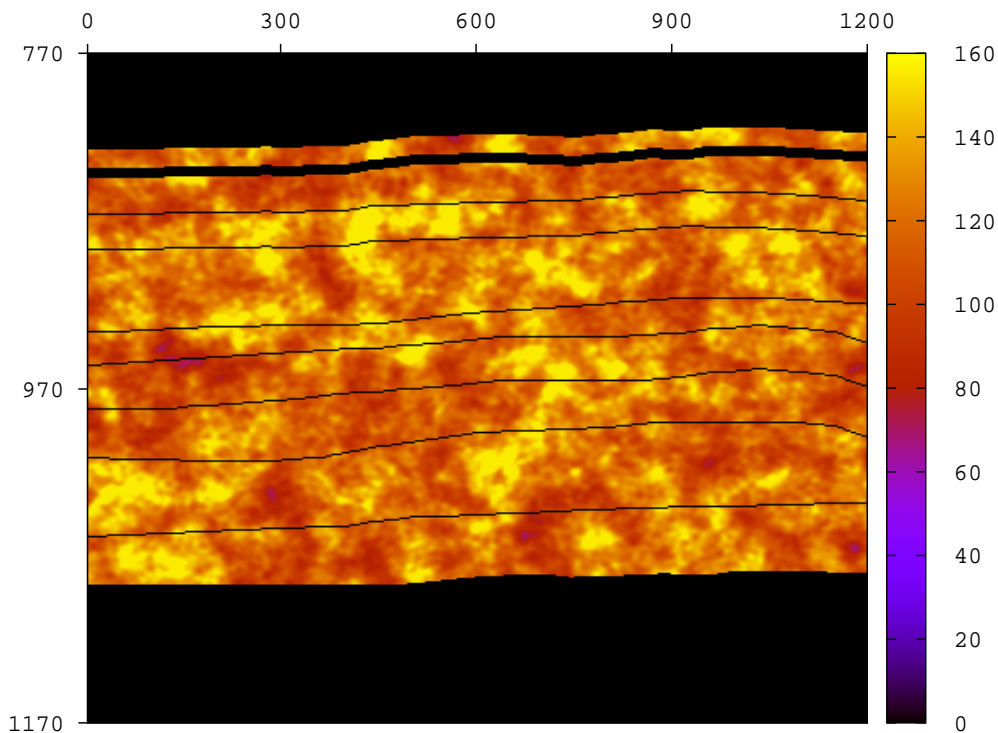


Figure 1: Vertical permeability map (mD)

The relative permeability curves adjusted from (Bachu and Bennion, 2008) with the power-law models (8, 9) are shown in Figure 2.

Besides, using the Leverett J -function we obtained the green $P_c(S_g)$ curve shown in Figure 3 for the mudstones and the red one for the Utsira formation (CASE 1).

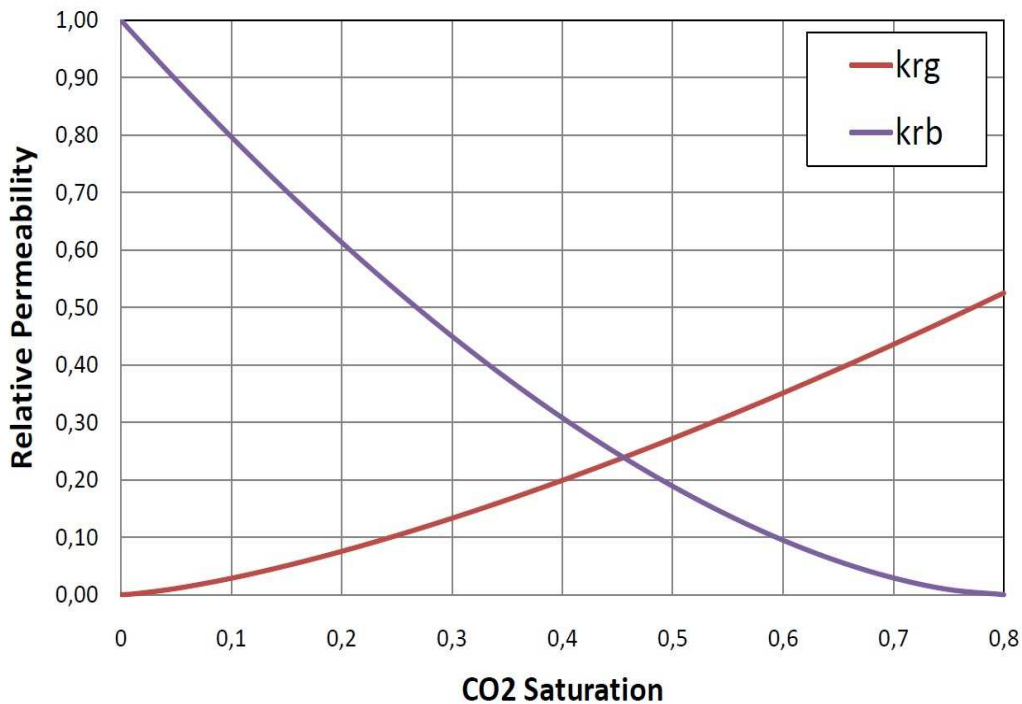


Figure 2: Relative Permeability Curves

To perform the sensitivity analysis, two other $P_c(S_g)$ curves for the Utsira formation, violet (CASE 2) and blue (CASE 3) in Figure 3, are used. They are obtained by multiplying the red curve values by 10 and 30, respectively.

5.2 CO₂ Injection

CO₂ is injected during three years in the Utsira formation at a constant flow rate of one million tons per year. The injection point is located at the bottom of the formation: $x = 0.6$ km, $y = 5.0$ km, $z = 1.082$ km. To satisfy the CFL stability condition due to IMPES formulation (Savioli and Bidner, 2005), the time step is 0.08 days. Recall that the petrophysical properties of the formation are time dependent due to the CO₂ injection and consequent increase in pore fluid pressure, but they change at a much slower rate than pressure and saturation. As a consequence, we have two time scales, and we use a much larger time step to update petrophysical properties than to run the flow simulator. In this work, the petrophysical properties are updated every 30 days.

Figure 4 shows the CO₂ saturation map after three years of injection obtained using the red (CASE 1) and green $P_c(S_g)$ curves in Figure 3 for the Utsira formation and mudstones, respectively. CO₂ accumulations below the mudstone layers can be observed. As CO₂ saturation increases, vertical permeability updated with equation 7 also increases, as it is shown in Figure 5. This allows a higher CO₂ mobility across the layers; therefore, part of the accumulated fluid migrates upwards, reaching the top of the formation. Porosity and horizontal permeability suffer little changes because they depend only on pressure.

Finally, we analyze the impact of capillary pressure curves on simulation results. Figures 6

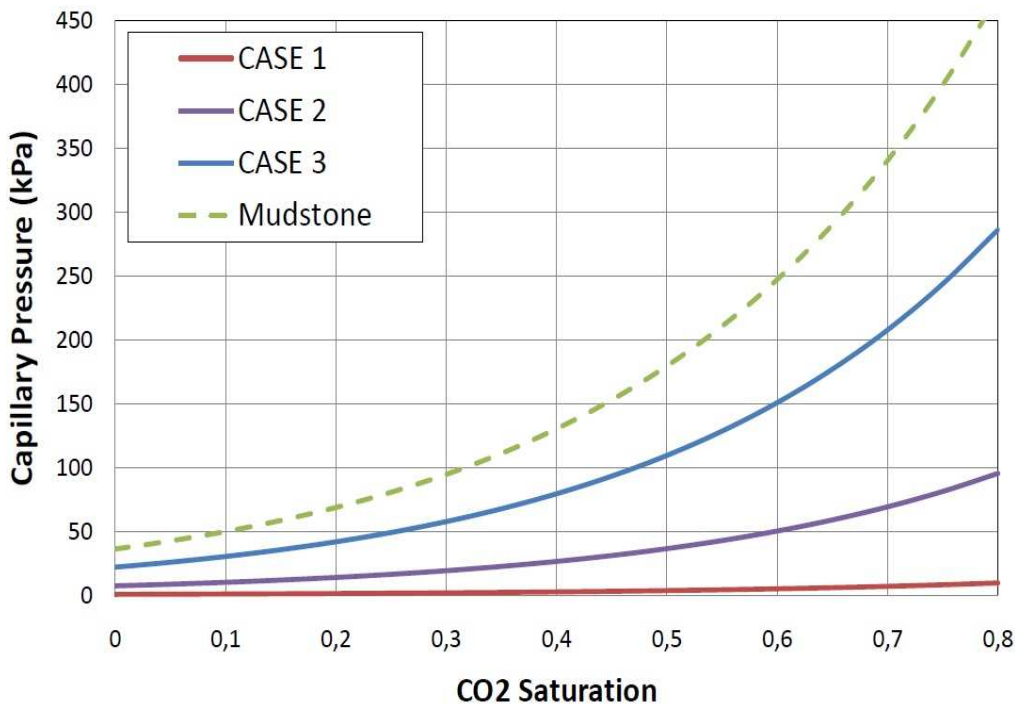


Figure 3: Capillary Pressure Curves (*kPa*)

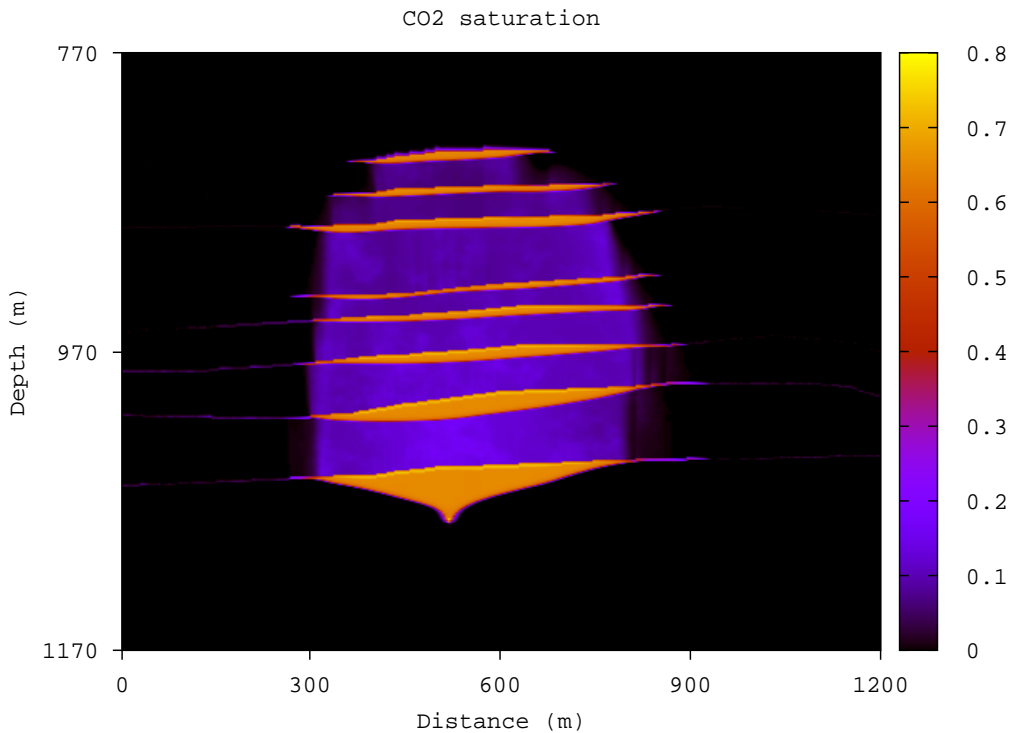


Figure 4: CO₂ Saturation map after 3 years of injection: CASE 1

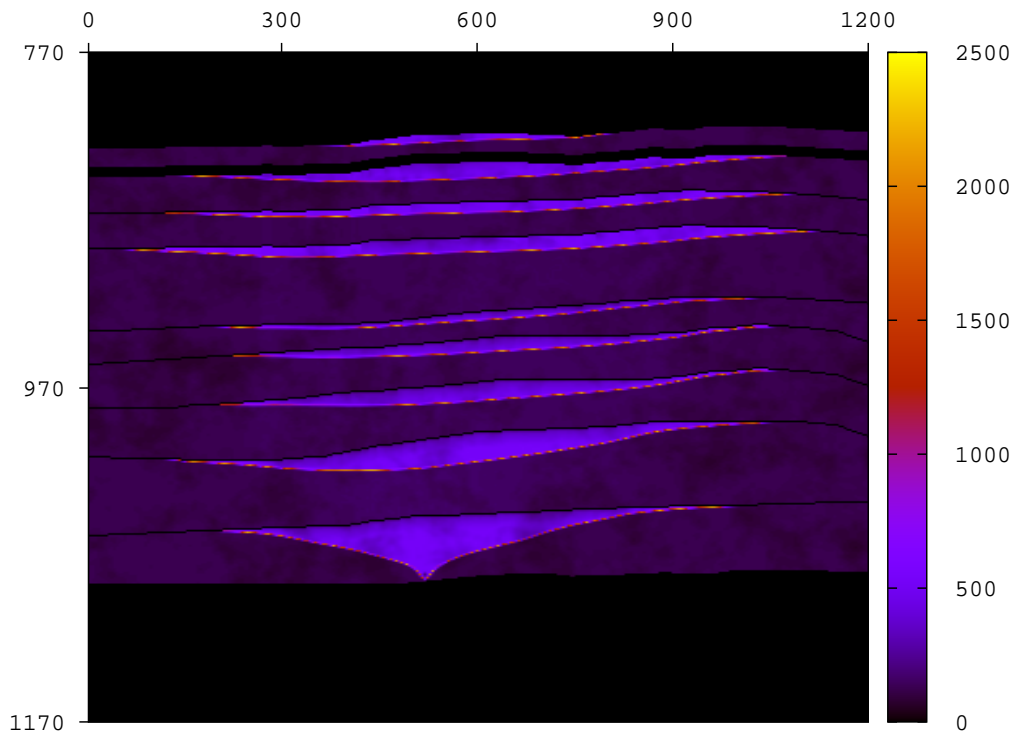


Figure 5: Vertical permeability map (mD) after three years: CASE 1

(CASE 2) and 7 (CASE 3) show the CO_2 saturation map after three years of injection obtained using the violet and blue capillary functions shown in figure 3.

Comparing Figures 4 (CASE 1), 6 (CASE 2) and 7 (CASE 3), we can observe that the increase in capillary pressure delays the CO_2 upward migration. As a consequence, there are higher CO_2 saturations levels between layers and accumulation zones spread laterally.

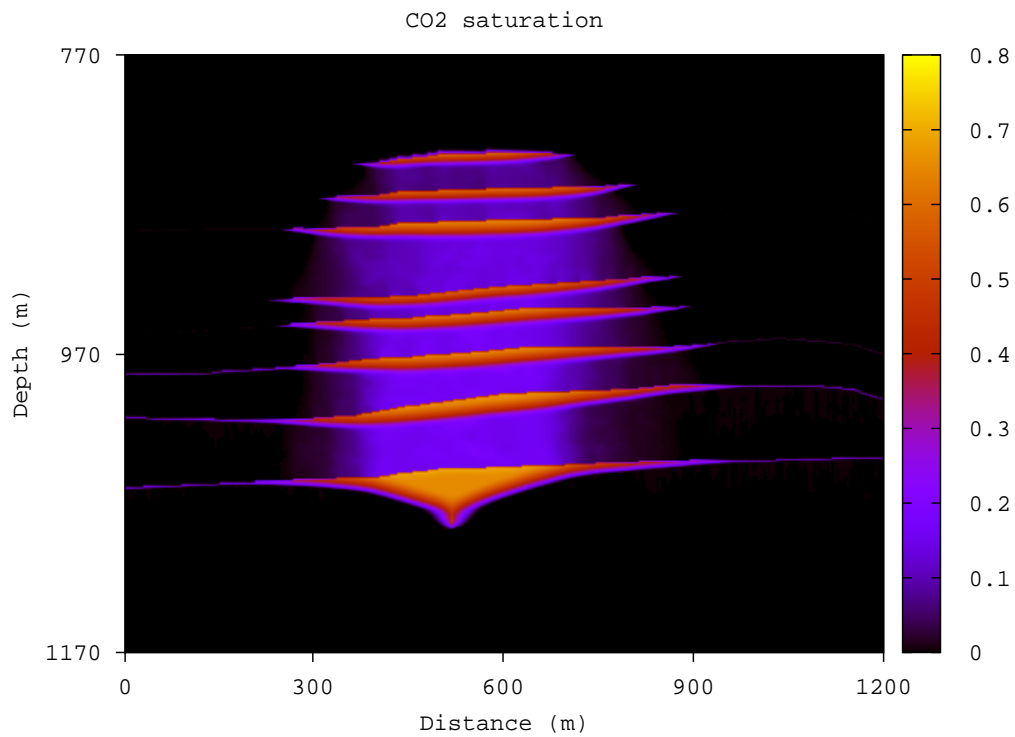
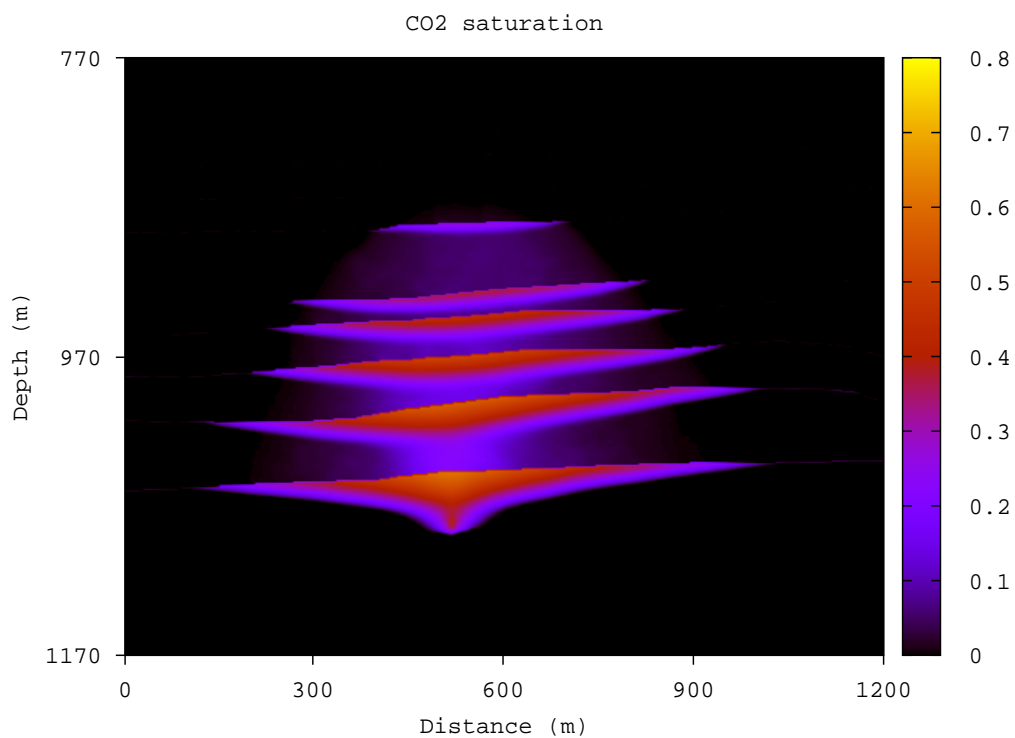
In addition, higher values of vertical permeability can be observed in CASE 2 (Figure 8) and CASE 3 (Figure 9), due to the higher CO_2 accumulation beneath the mudstone layers. However, this vertical permeability increment that facilitates upward motion, is not enough to compensate the capillary pressure effect.

6 CONCLUSIONS

The fluid-flow simulator yields CO_2 accumulations below the mudstone layers. Capillary forces affect the migration and dispersal of the CO_2 plume; higher values of these forces cause a slower CO_2 upward migration and lateral spreading of accumulation zones below mudstone layers.

Therefore, the influence of capillary pressure is very important in CO_2 storage; consequently an accurate estimation of this function is essential to perform realistic long term predictions.

Numerical simulation constitutes a meaningful tool to analyze CO_2 storage integrity, helping to prevent any leakage back up to the atmosphere or into underground sources of drinking water.

Figure 6: CO₂ Saturation map after 3 years of inyection: CASE 2Figure 7: CO₂ Saturation map after 3 years of inyection: CASE 3

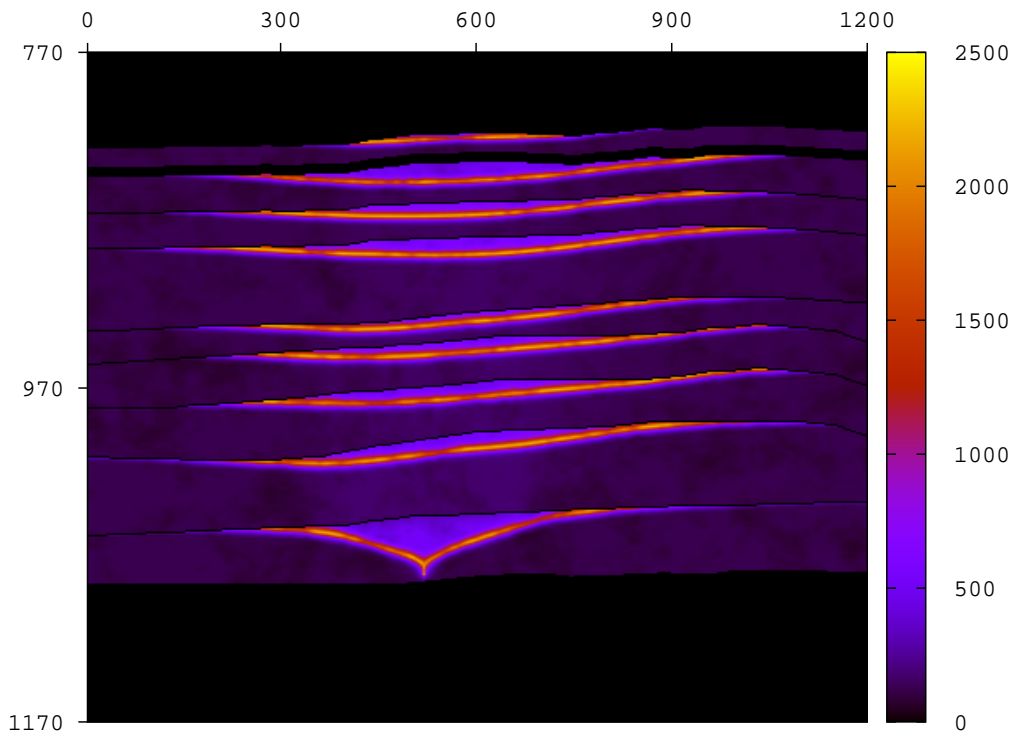


Figure 8: Vertical permeability map (mD) after three years: CASE 2

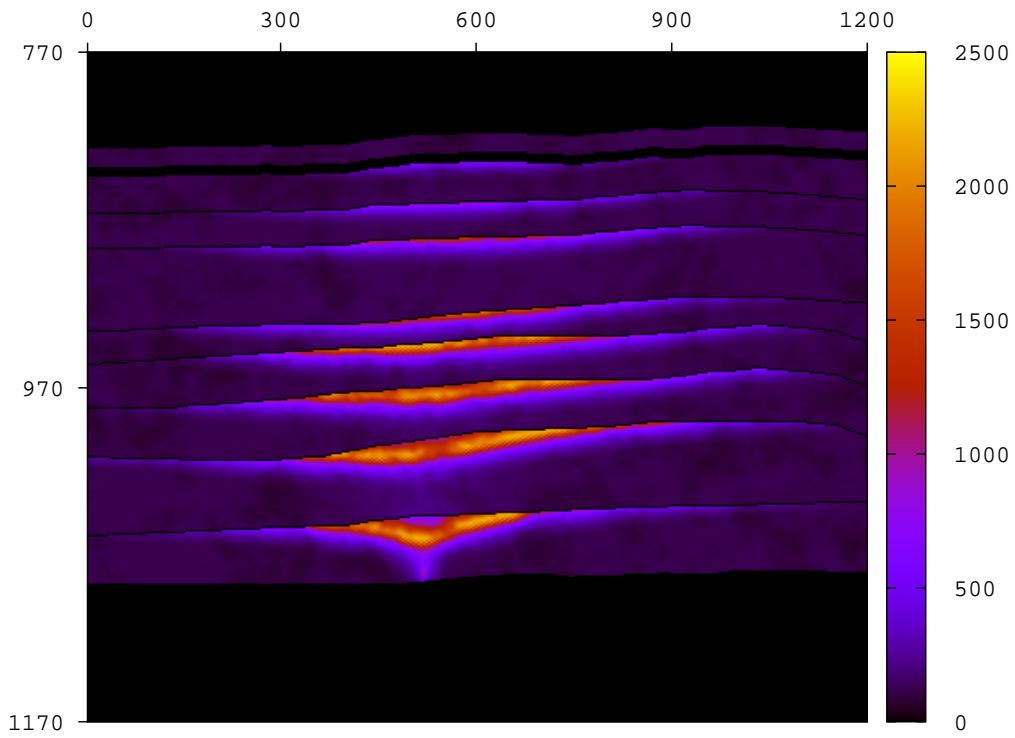


Figure 9: Vertical permeability map (mD) after three years: CASE 3

7 ACKNOWLEDGEMENTS

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