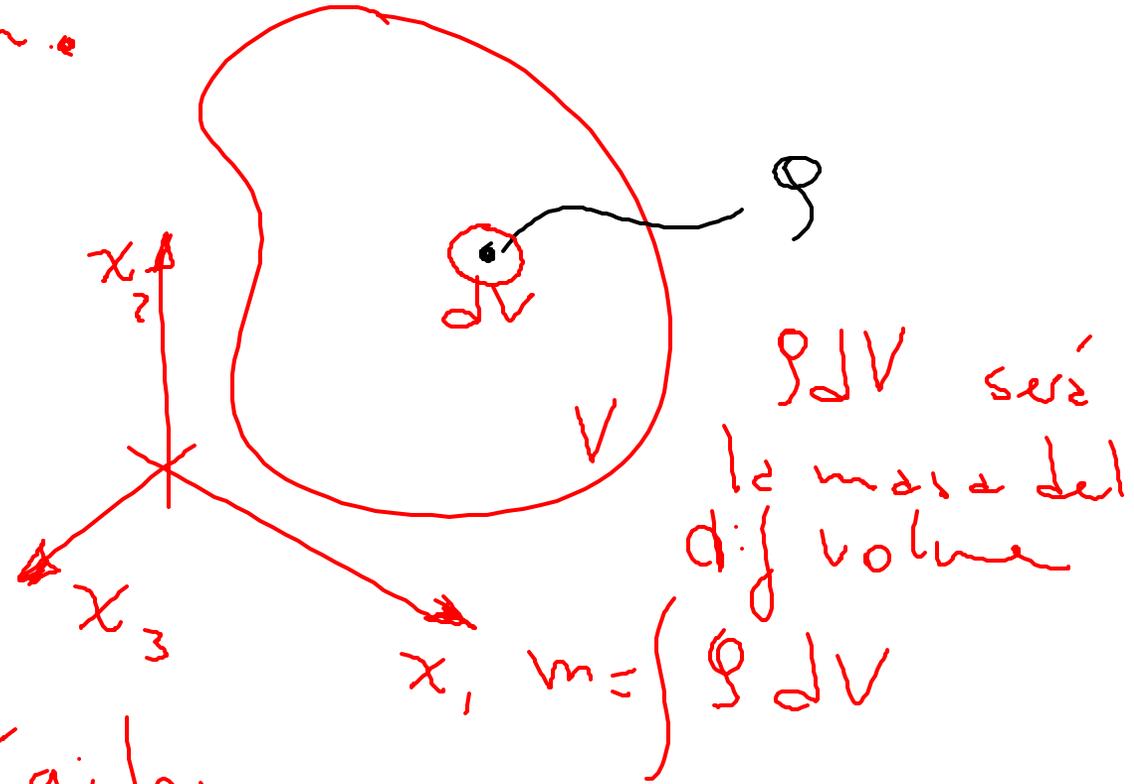


TEMAS: Mov cuerpo rígido

CR: sistema de partículas / las distancias entre partículas no varían.

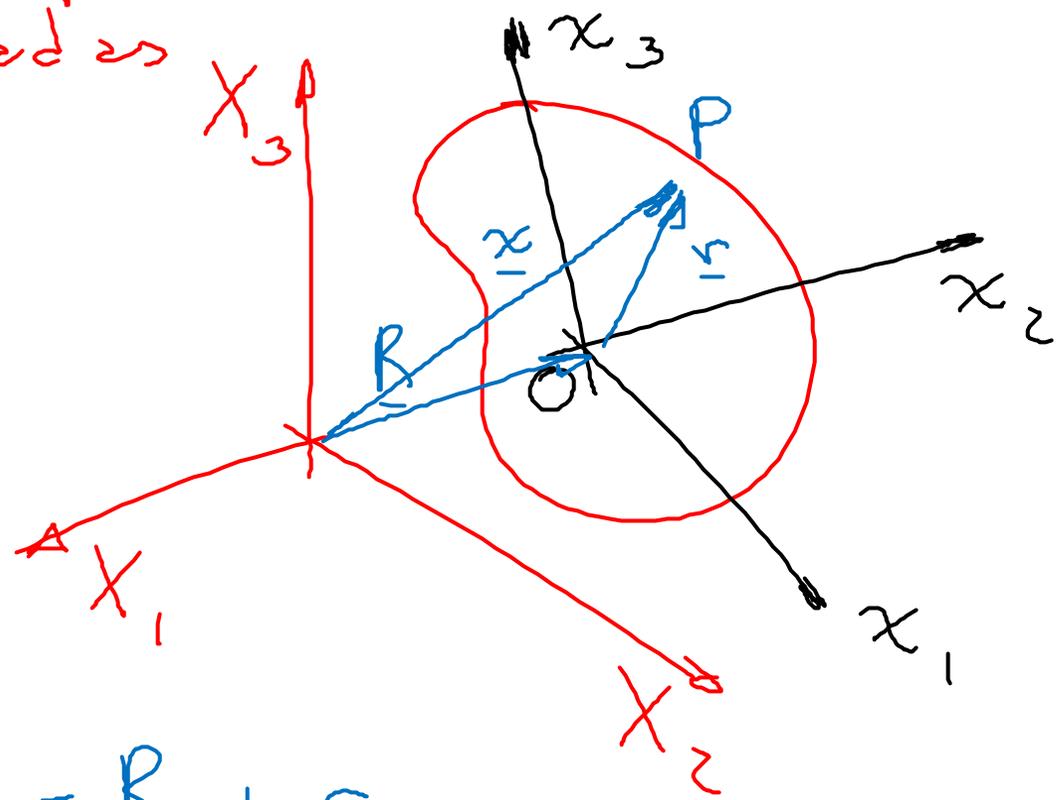
Es una aproximación.

Los cuerpos sólidos muchas veces verifican esta condición.



Cuerpos sólidos { Rígidos
Flexibles { Elasticidad
Plasticidad
V P last

P/descripción del movimiento, usamos dos sistemas de coordenadas



X_1, X_2, X_3 : sistema "fijo" o inercial
"Coord espaciales"

$O x_1, x_2, x_3$: sistema fijo en el cuerpo
"Coord materiales"

O : fijo en el centro de masa del cuerpo.

$$\underline{x} = \underline{R} + \underline{r}$$

P/determinar la posición de P , necesitamos conocer coord del CM (\underline{R}) y la orientación del (Ox_1, x_2, x_3) .

La orientación del cuerpo la podemos determinar, por ejemplo, conociendo 3 ángulos de Euler (u otro sist)

luego, el CR tiene 6 GDL.

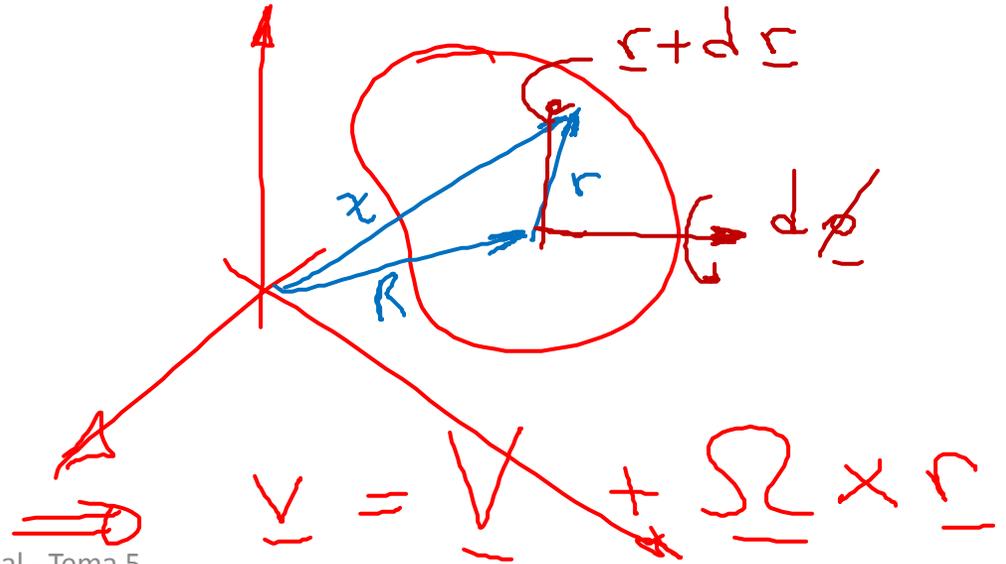
Consideremos un elemento infinitesimal arbitrario del CR. Lo podemos ver como P/2 partes:

- i) traslación infinitesimal y orientación fija
- ii) rotación infinitesimal entorno al centro de masa.

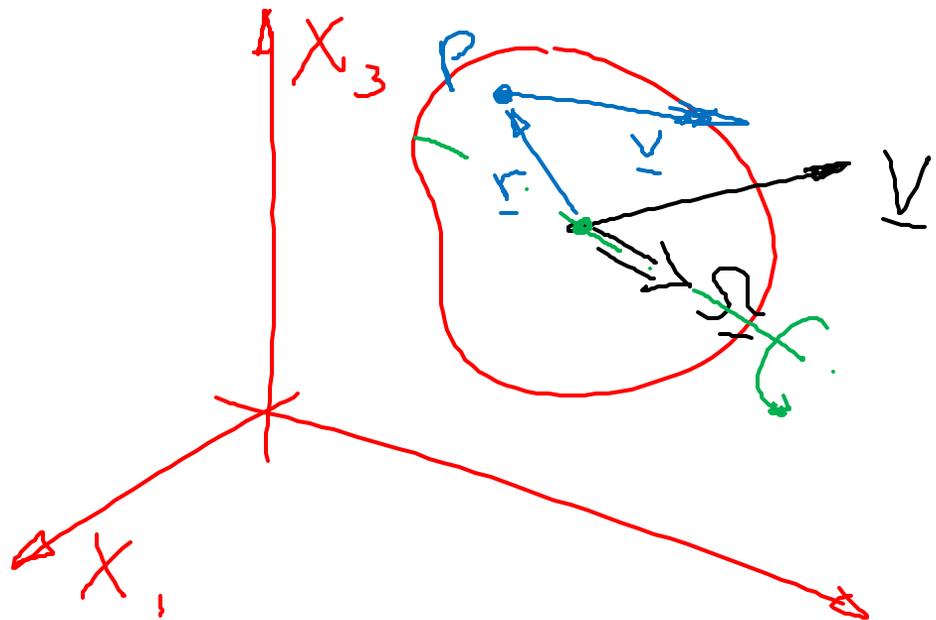
$$\underline{x} = \underline{R} + \underline{r}$$

$$d\underline{x} = d\underline{R} + d\phi \times \underline{r}$$

$$\frac{d\underline{x}}{dt} = \frac{d\underline{R}}{dt} + \frac{d\phi}{dt} \times \underline{r}$$



$D_{iv} \times dt:$



$$\underline{v} = \underline{V} + \underline{\Omega} \times \underline{r}$$

\underline{V} : vel de traslación del CR

$\underline{\Omega}$: vel angular de rotación del CR

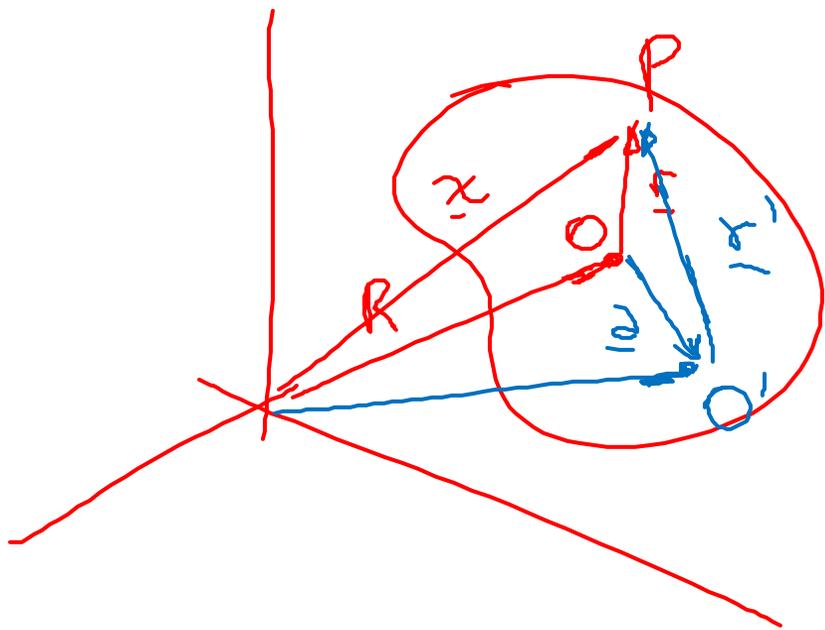
Assumimos origen de coord del sist material esté en un punto O' a una dist \underline{a} del punto O (Centro de masa).

Sea \underline{V}' la vel de O'

$\underline{\Omega}'$ la vel angular del nuevo sist

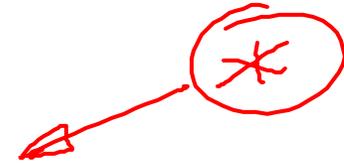
\underline{r}' la posición de P en el nuevo sist

$$\underline{r} = \underline{r}' + \underline{a}$$



$$\underline{r} = \underline{\rho} + \underline{r}'$$

En la anterior:



$$\begin{aligned} \underline{v} &= \underline{V} + \underline{\Omega} \times \underline{r} = \\ &= \underline{V} + \underline{\Omega} \times \underline{\rho} + \underline{\Omega} \times \underline{r}' \\ &\quad \underbrace{\hspace{10em}}_{\underline{V}'} \end{aligned}$$

Entonces:

$$\underline{v} = \underline{V}' + \underline{\Omega} \times \underline{r}'$$

y por comparación con *

concluyo que

$$\underline{\Omega}' = \underline{\Omega}$$

Todo punto del CA
está animado de la misma
vel de rotación.

Luego, hablo de "velocidad angular del cuerpo rígido".

De la expresión:

$$\underline{V}' = \underline{V} + \underline{\Omega} \times \underline{a}$$

venos q/si $\underline{V} \perp \underline{\Omega}$, luego \underline{V}' y $\underline{\Omega}'$ son perpendiculares

Por cualquier elección de O' .

$$\left(\underline{V}' = \underline{V} + (\underline{V}' - \underline{V}) \quad \text{y} \quad (\underline{V}' - \underline{V}) = \underline{\Omega} \times \underline{a} \perp \underline{\Omega} \right)$$

Luego, es posible encontrar \underline{a} / $\underline{V}' = \underline{0}$

$$\underline{V} + \underline{\Omega} \times \underline{a} = \underline{0}$$

$$\underline{V} + \underline{\Omega} \times \underline{a} = \underline{0} \quad \Rightarrow \quad \underline{\tilde{\Omega}} \underline{a} = -\underline{V}$$

$$\underline{\tilde{\Omega}} = \begin{bmatrix} 0 & \Omega_3 & -\Omega_2 \\ -\Omega_3 & 0 & \Omega_1 \\ \Omega_2 & -\Omega_1 & 0 \end{bmatrix}$$

$\underline{\tilde{\Omega}}$: matriz antisimétrica / $\forall \underline{a}$
 $\underline{\tilde{\Omega}} \underline{a} = \underline{\Omega} \times \underline{a}$

$$\underline{V} \cdot \underline{\Omega} = 0 \quad \Rightarrow \quad \underline{V}' \cdot \underline{\Omega}' = 0$$

Si \underline{a} es solución de (*), luego $(\underline{a} + \alpha \underline{\Omega})$ también lo es:

$$\underline{\tilde{\Omega}} (\underline{a} + \alpha \underline{\Omega}) = \underline{\tilde{\Omega}} \underline{a} + \alpha \underbrace{\underline{\tilde{\Omega}} \underline{\Omega}}_0 = -\underline{V}$$

Sea $\alpha / \underline{a}^* \cdot \underline{\Omega} = 0$

$$\begin{bmatrix} \tilde{\Omega} \\ \Omega \end{bmatrix} \underline{a} = \begin{pmatrix} -V \\ 0 \end{pmatrix}$$

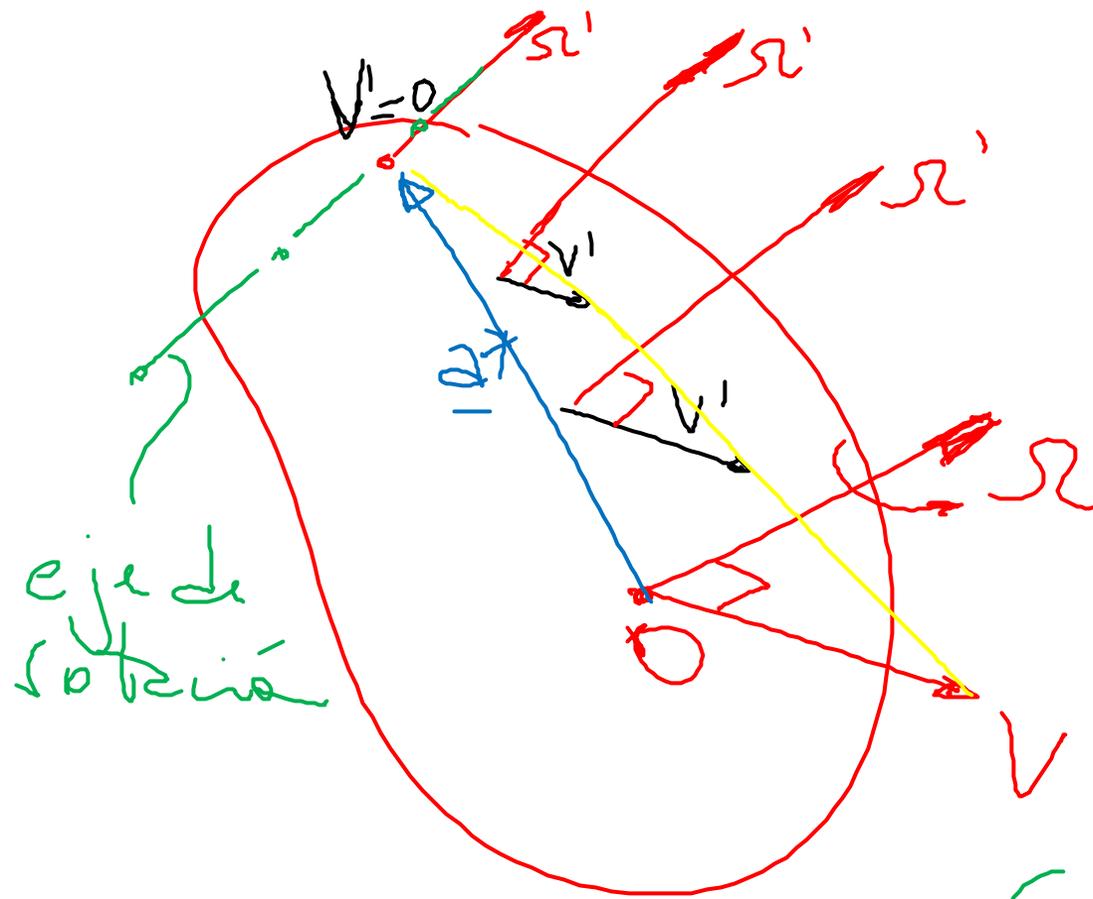
$$\underline{\tilde{\Omega}}^2 = \underline{\Omega} \otimes \underline{\Omega} - \Omega^2 \underline{I}$$

0	Ω_3	$-\Omega_2$	$-\Omega_3^2 - \Omega_2^2$	$\Omega_1 \Omega_2$
$-\Omega_3$	0	Ω_1	$\Omega_1 \Omega_2$	$-\Omega_3^2 - \Omega_1^2$
Ω_2	$-\Omega_1$	0	$\Omega_1 \Omega_3$	

$-\tilde{\Omega}$	Ω	$\Omega^2 \underline{I} - \Omega \otimes \Omega + \Omega \times \Omega$
-------------------	----------	---

$$\underline{a} = \frac{1}{\Omega^2} \begin{bmatrix} -\tilde{\Omega} & \Omega \end{bmatrix} \begin{pmatrix} -V \\ 0 \end{pmatrix} =$$

$$\underline{a} = \frac{\Omega \times V}{\Omega^2}$$



i) Todo $\underline{v}' \perp \underline{\Omega}$

$$\underline{v}^* = \frac{\underline{\Omega} \times \underline{v}}{\Omega^2}$$

$$\underline{v} = \underline{v}^* + \alpha \underline{\Omega}$$

α arbitrario

Cualquier pto sobre el eje de rotación tiene velocidad $\underline{v}' = \underline{0}$

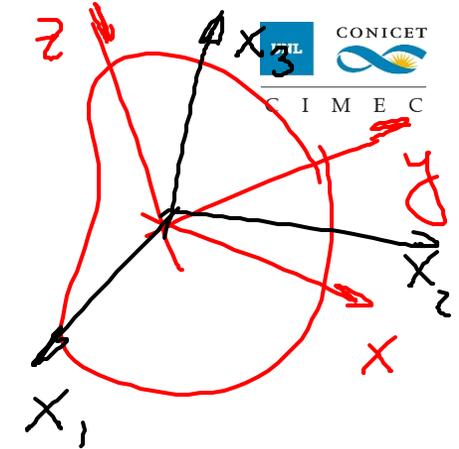
$$\begin{aligned}
 T_{rot} &= \frac{1}{2} \sum m \left(\Omega^2 r^2 - (\underline{\Omega} \cdot \underline{r})^2 \right) = \\
 &= \frac{1}{2} \underline{\Omega} \cdot \sum m r^2 \underline{I} \underline{\Omega} - \frac{1}{2} \underline{\Omega} \cdot \sum m \underline{r} \otimes \underline{r} \underline{\Omega} \\
 &= \frac{1}{2} \underline{\Omega} \cdot \underbrace{\left[\sum m r^2 \underline{I} - m \underline{r} \otimes \underline{r} \right]}_{\underline{J}} \underline{\Omega}
 \end{aligned}$$

\underline{J} tensor de inercia del CR

$$T = \frac{1}{2} m V^2 + \frac{1}{2} \underline{\Omega} \cdot \underline{J} \underline{\Omega}$$

\underline{J} es un tensor simétrico

$$J = \begin{bmatrix} \sum m (y^2 + z^2) & -\sum m y x & -\sum m z x \\ -\sum m y x & \sum m (x^2 + z^2) & -\sum m z y \\ -\sum m z x & -\sum m z y & \sum m (x^2 + y^2) \end{bmatrix} \quad \text{SIM}$$



J_{xx}, J_{yy}, J_{zz} momentos de inercia resp de ejes x, y, z

TENSOR ADITIVO

Como todo tensor, puede ser diagonalizado.

En los ejes $x_1, x_2, x_3 \Rightarrow$
"Ejes ppales de inercia"

$$J = \begin{bmatrix} J_1 & 0 & 0 \\ 0 & J_2 & 0 \\ 0 & 0 & J_3 \end{bmatrix}$$

En los ejes ppales de inercia:

$$T_{\text{rot}} = \frac{1}{2} \left(I_1 \Omega_1^2 + I_2 \Omega_2^2 + I_3 \Omega_3^2 \right) = \frac{1}{2} \underline{\Omega} \cdot \underline{I} \underline{\Omega}$$

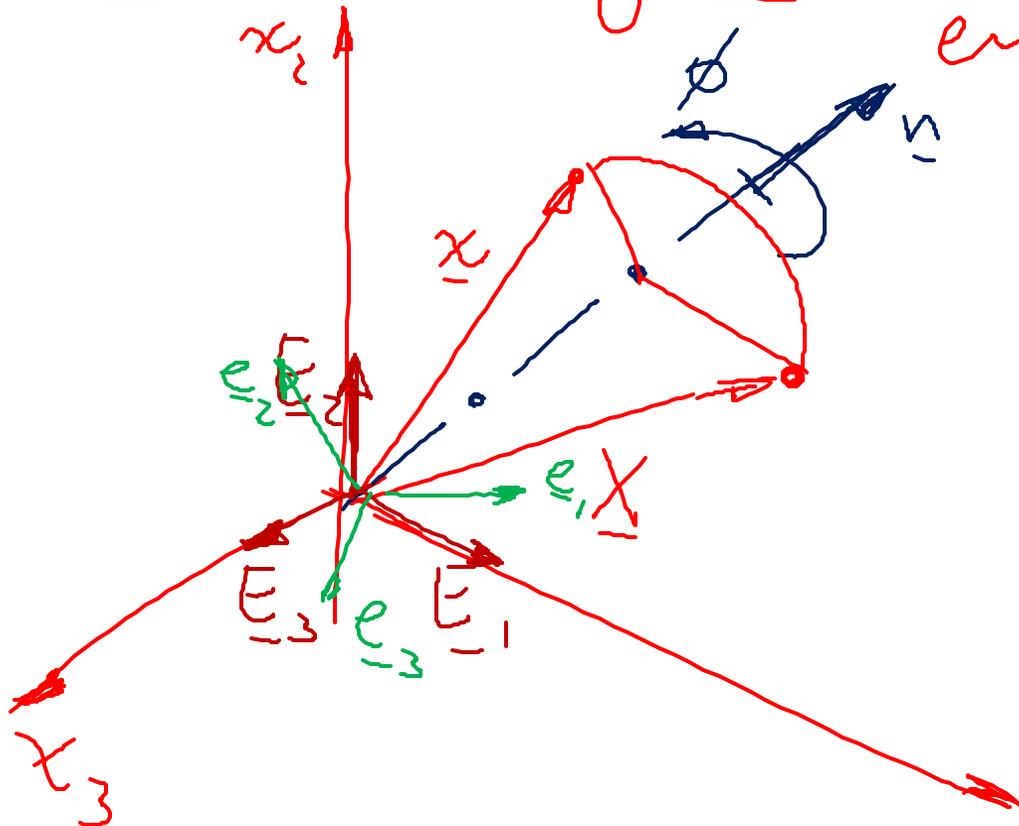
CR no tiene simetrías \Rightarrow trompo asimétrica.

CR \sphericalangle 2 momentos ppals iguales $(I_1 = I_2) \Rightarrow$ " simétrica.

CR \sphericalangle 3 momentos ppals iguales $(I_1 = I_2 = I_3) \Rightarrow$ " esférico

CINEMÁTICA DE ROTACIONES FINITAS

Movimiento esférico : corresponde a la rotación de un CR
entorno a un punto fijo en el espacio



$$\underline{X} = \begin{pmatrix} X_1 \\ X_2 \\ X_3 \end{pmatrix}$$

posición de P
en la ref. referencial

$$\underline{x} = \begin{pmatrix} x_1 \\ x_2 \\ x_3 \end{pmatrix}$$

posición actual
de P.

$\underline{E}_1, \underline{E}_2, \underline{E}_3$ vect. ortomomiales
fijos al cuerpo.

$\underline{e}_1, \underline{e}_2, \underline{e}_3$ los mismos
vect. ortomomiales
luego de la transformación

Mov CR

- i) la long del vector posición de P debe ser cte
- ii) el áng relativo entre 2 vect arbitrarios fijos al cuerpo debe ser cte.

La transf de $X \rightarrow x$ puede ser expresada por una transf lineal

$$\underline{x} = \underline{R} \underline{X}$$

de la misma manera

$$\underline{e}_i = \underline{R} \underline{E}_i$$

Como la longitud del radio vector es cte:

$$\underline{x}^T \underline{x} = \underline{X}^T \underline{X} = \underline{X}^T \underline{R}^T \underline{R} \underline{X} \Rightarrow \underline{R}^T \underline{R} = \underline{I}$$

$$\underline{R} \text{ es ortogonal} \Rightarrow \underline{R}^T = \underline{R}^{-1}$$

llamado

$$\underline{A} = [\underline{E}_1 \ \underline{E}_2 \ \underline{E}_3] \quad \underline{B} = [\underline{e}_1 \ \underline{e}_2 \ \underline{e}_3]$$

$$\underline{B} = \underline{R} \underline{A}$$

Siendo \underline{E}_i unitarios \Rightarrow
 como dextrógira

$$\det \underline{A} = 1$$

$$\det \underline{B} = 1$$

$\Rightarrow \det \underline{R} = 1$
 \underline{R} : matriz ortog.
 propia

Autovectores de $\underline{\underline{R}}$

$$\det \underline{\underline{R}} = \lambda_1 \lambda_2 \lambda_3 = 1$$

Se puede verificar que λ_i tiene la forma:

$$\lambda_1 = 1$$

$$\lambda_{2,3} = \exp(\pm i\phi) = \cos \phi \pm i \sin \phi$$

Luego, $\underline{\underline{R}}$ admite al menos un autovector \underline{n} /

$$\underline{\underline{R}} \underline{n} = \underline{n} \quad \underline{n} : \text{eje de rotación!}$$

Los autovectores asociados a $\lambda_{2,3}$ serán complejos conjugados

Expresos los autovectores de \underline{R} en una matriz \underline{X} :

$$\underline{X} = \begin{bmatrix} n & \underline{u} + i\underline{v} & \underline{u} - i\underline{v} \end{bmatrix}$$

El prob de autovalores, en forma matricial, se expresa:

$$\underline{R} \underline{X} = \underline{X} \underline{\Lambda}$$

$$\underline{R} \underline{X} \underline{X}^* = \underline{R} = \underline{X} \underline{\Lambda} \underline{X}^*$$

$$\underline{X}^{-1} = \underline{X}^* \quad (\text{transp conjugado})$$

$$X^* = \begin{bmatrix} \underline{n}^T \\ (\underline{u} - i\underline{v})^T \\ (\underline{u} + i\underline{v})^T \end{bmatrix}$$

$$X^* X = \begin{bmatrix} \underline{n}^T \underline{n} & \underline{n}^T (\underline{u} + i\underline{v}) \\ (\underline{u} - i\underline{v})^T (\underline{u} + i\underline{v}) & \vdots \\ \vdots & \vdots \end{bmatrix} = \underline{I}$$

$$\underline{n}^T \underline{u} = \underline{n}^T \underline{v} = 0$$

$$\underline{u}^T \underline{v} = 0$$

$$\underline{u}^T \underline{u} + \underline{v}^T \underline{v} = 1$$

$\underline{u}, \underline{v}$ forma una ~~base~~ ~~plano~~ ~~de~~ ~~vector~~ ~~ortog~~ ~~en~~ ~~un~~ ~~plano~~ ~~normal~~ ~~a~~ \underline{n}

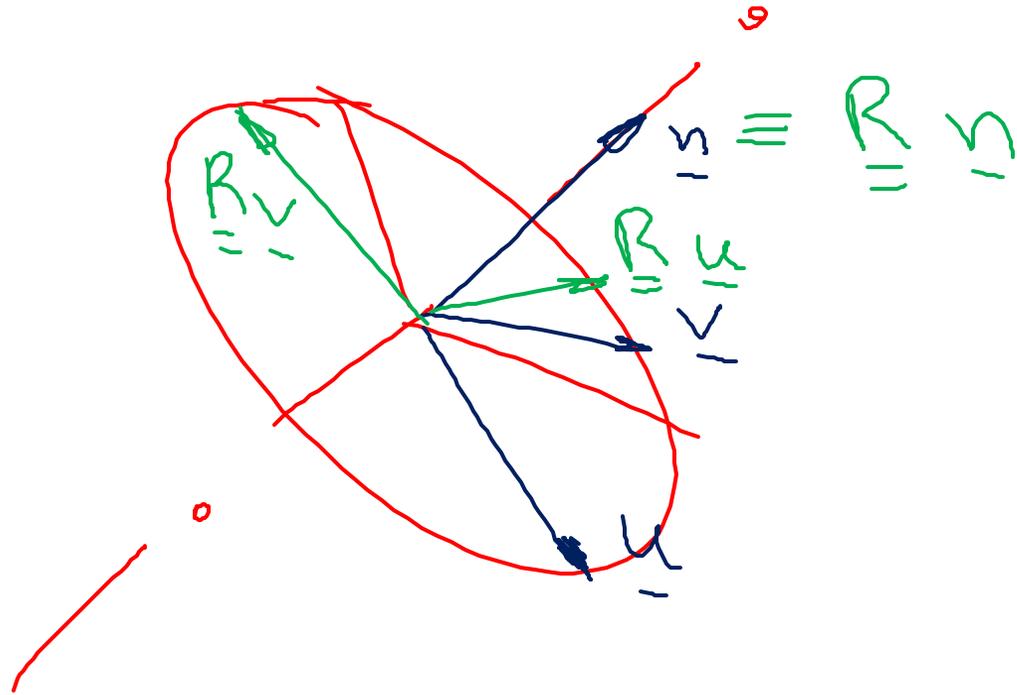
$$\underline{R}(\underline{u} + i\underline{v}) = \exp(i\phi) (\underline{u} + i\underline{v})$$

Desarrollando:

$$\underline{R}\underline{u} = \underline{u} \cos\phi - \underline{v} \sin\phi$$

$$\underline{R}\underline{v} = \underline{u} \sin\phi + \underline{v} \cos\phi$$

○ sea, $\forall \underline{u}, \underline{v}$ sufrir una rotación plana de ángulo ϕ en un plano perpendicular a \underline{n} .

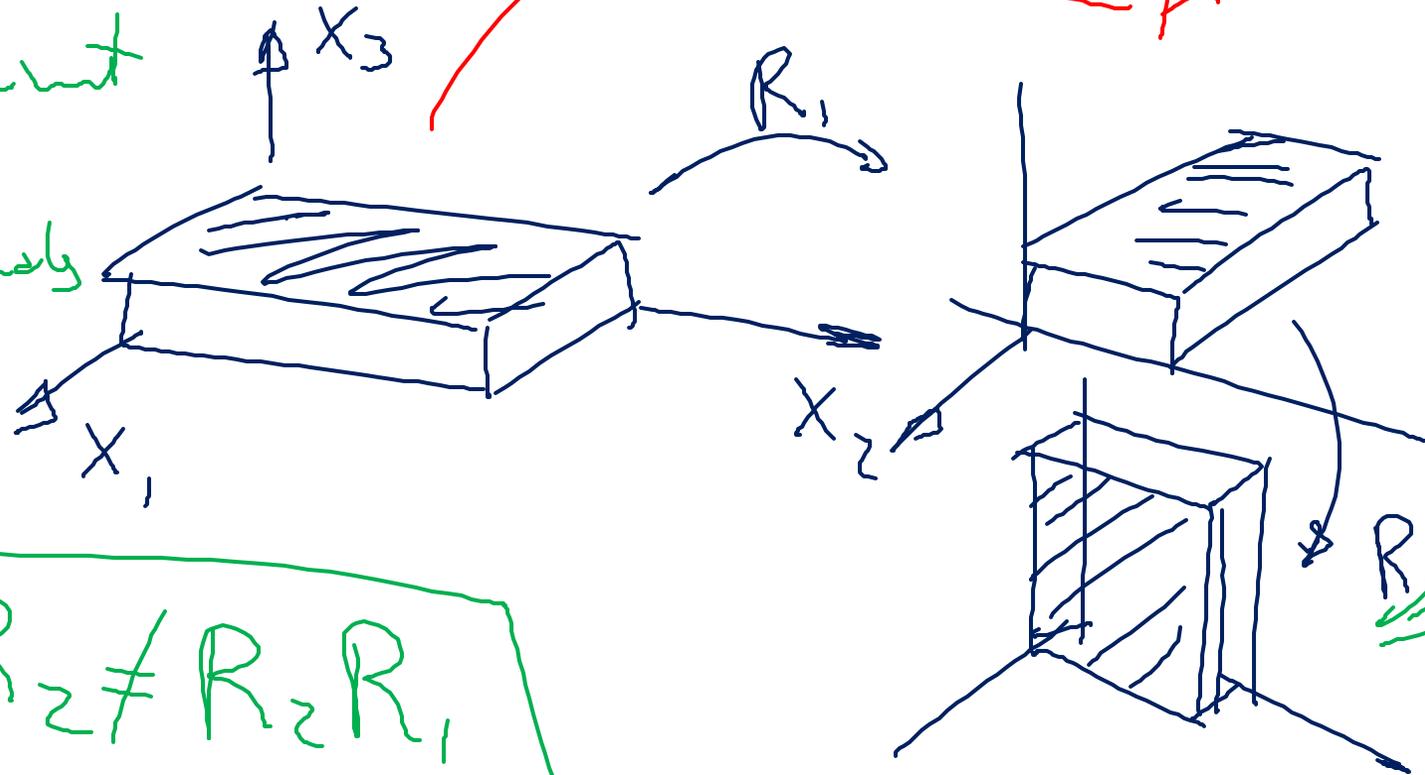


Rotaciones son no conmutativas

$$R_1 = R(z, 90^\circ)$$

$$R_2 = R(y, 90^\circ)$$

Se verifica con un
en el caso de
rot infinitesimales



$$R_1 R_2 \neq R_2 R_1$$

Expresión del tensor de rotación

$$\underline{\underline{R}}^T \underline{\underline{R}} = \underline{\underline{I}}$$

 \implies

6 ecu de restricción

$$\underline{\underline{R}} = [\underline{r}_1 \ \underline{r}_2 \ \underline{r}_3]$$

$$\implies r_i^T r_j = \delta_{ij}$$

Luego

$$\underline{\underline{R}} = \underline{\underline{R}}(\alpha_1, \alpha_2, \alpha_3)$$

función de 3
parámetros

Operador de proyección

$$\underline{\underline{P}}_n = \underline{\underline{n}} \underline{\underline{n}}^T \underline{\underline{n}} = \underline{\underline{I}} - \underline{\underline{n}} \underline{\underline{n}}^T$$

$$\underline{\underline{P}}_n \underline{\underline{v}} = \underline{\underline{v}} - \underline{\underline{n}} \left(\underline{\underline{n}}^T \underline{\underline{v}} \right) =$$

Tensor spin

$$\underline{\underline{S}}_n = \begin{bmatrix} 0 & -n_3 & n_2 \\ n_3 & 0 & -n_1 \\ -n_2 & n_1 & 0 \end{bmatrix} = \text{spin}(\underline{\underline{n}})$$

$$\underline{\underline{S}}_n \underline{\underline{v}} = \underline{\underline{n}} \times \underline{\underline{v}}$$

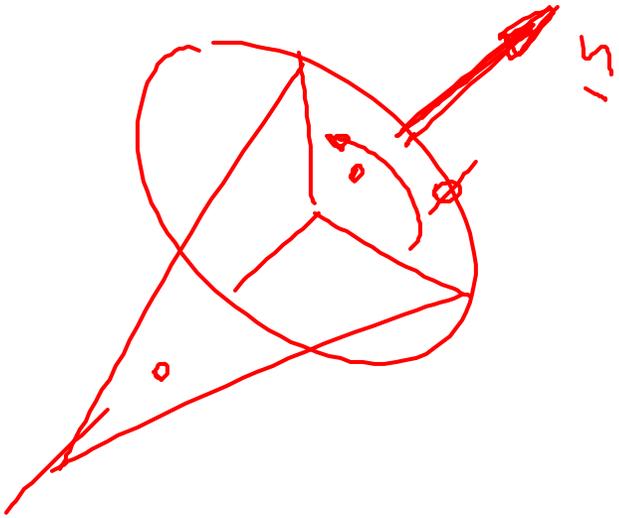
$$\underline{\underline{S}}_n \underline{\underline{n}} = \underline{\underline{0}}$$



$$\underline{\underline{n}}^T \underline{\underline{n}} = 1$$

			n_3	0
			$-n_2$	n_1
0	n_3	$-n_2$	$n_3^2 + n_2^2$	
$-n_3$	0	n_1	$-n_1 n_2$	$n_3^2 + n_1^2$
n_2	$-n_1$	0	$-n_1 n_3$	

$$\underline{R} = \underline{I} \cos \phi + (1 - \cos \phi) \underline{n} \underline{n}^T + \underline{\hat{n}} \sin \phi \quad \Bigg| \quad \left(= \exp(\underline{\hat{n}} \phi) \right)$$



$$\underline{x} = \underline{R} \underline{X}$$

\underline{n}, ϕ
 $\text{pero } \underline{n}^T \underline{n} = 1$
 3 parámetros libres

Mapeo exponencial

$$\underline{x} = \underline{R} \underline{X}$$

Diferenciamos respecto de ϕ

$$\frac{d\underline{x}}{d\phi} = \frac{d\underline{R}}{d\phi} \underline{X} = \frac{d\underline{R}}{d\phi} \underline{R}^T \underline{x}$$

$$\frac{d\underline{R}}{d\phi} = \underline{\Omega} \cos\phi - (\underline{I} - \underline{nn}^T) \sin\phi$$

Se verifica fácilmente que $\frac{d\underline{R}}{d\phi} \underline{R}^T = \underline{\Omega}$

EJERCICIO

Luego

$$\begin{cases} \frac{d\underline{x}}{d\phi} - \underline{\hat{n}} \underline{x} = 0 & \text{Ec dif. } \phi \\ \underline{x}(0) = \underline{X} \end{cases}$$

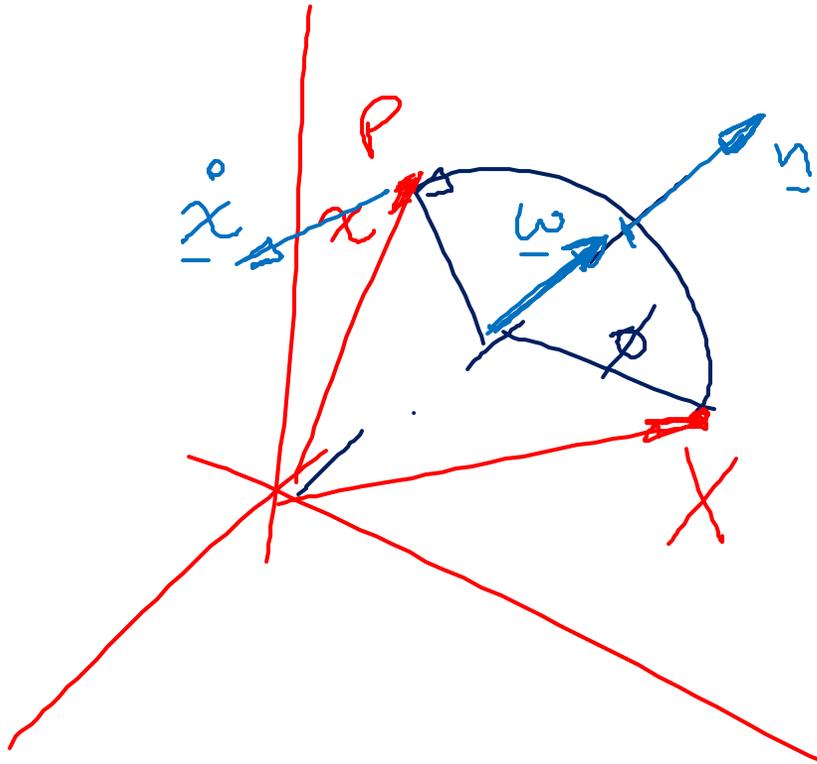
La solución de este EDO

$$\underline{x} = \exp(\underline{\hat{n}}\phi) \underline{X}$$

o o

$$\underline{R} = \exp(\underline{\hat{n}}\phi) \underline{I}$$

Análisis de velocidad P/ mov esférico



$$x = R X \quad R^T = R^{-1}$$

$$v_p = \dot{x} = \dot{R} X = \dot{R} R^T x$$

$$\frac{d}{dt} (R R^T) = \dot{R} R^T + R \dot{R}^T = 0$$

$$\dot{R} R^T + (\dot{R} R^T)^T = 0$$

$\dot{R} R^T$ antisimétrica

matriz de velocidades angulares expresada en coord espaciales

Definimos $\tilde{\omega} = \dot{R} R^T$

$$\dot{x} = \tilde{\omega} x = \omega \times x$$

$$\underline{\omega} = \text{vect}(\underline{\dot{R}}\underline{R}^T)$$

La vel del punto P la
puedo expresar en coord
Materiales.

$$V_P = R^T v_P = \underbrace{R^T \dot{R}}_{AS} X$$

Definición:

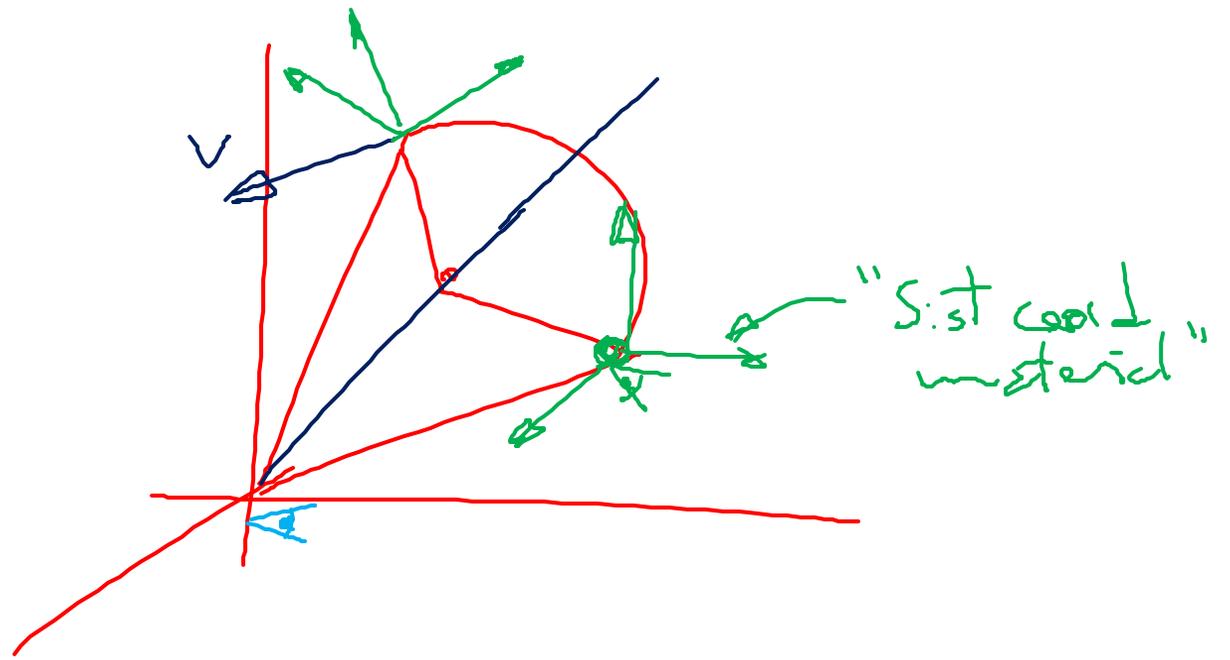
$$\underline{\underline{\Omega}} = \underline{\underline{R}}^T \dot{\underline{\underline{R}}}$$

$$\underline{\underline{\Omega}} = \text{vect}(\underline{\underline{R}}^T \dot{\underline{\underline{R}}})$$

"matriz de velocidades angulares materiales"

$$\dot{\underline{\underline{R}}} = \underline{\underline{\Omega}} \underline{\underline{R}} = \underline{\underline{R}} \underline{\underline{\Omega}}$$

$$\underline{\underline{\Omega}} = \underline{\underline{R}} \underline{\underline{\Omega}} \underline{\underline{R}}^{-1} \quad \underline{\underline{\Omega}} = \underline{\underline{R}}^{-1} \dot{\underline{\underline{R}}} \underline{\underline{R}}$$



$$\underline{\tilde{\omega}} = \underline{R} \underline{\tilde{\Omega}} \underline{R}^T = \underline{\tilde{(R\Omega)}}$$

$$\underline{\omega} = \underline{R} \underline{\Omega} \implies \underline{\Omega} = \underline{R}^T \underline{\omega}$$

$$\begin{aligned} \underline{\tilde{\omega}} \cdot \underline{v} &= \underline{R} \underline{\tilde{\Omega}} (\underline{R}^T \underline{v}) = \underline{R} (\underline{\Omega} \times \underline{R}^T \underline{v}) = \\ &= \underline{R} \underline{\Omega} \times (\underline{R} \underline{R}^T \underline{v}) = (\underline{R} \underline{\Omega}) \times \underline{v} = \underline{\tilde{(R\Omega)}} \cdot \underline{v} \end{aligned}$$

\underline{v} arbitrario $\implies \underline{\tilde{\omega}} = \underline{\tilde{(R\Omega)}} \implies \underline{\omega} = \underline{R} \underline{\Omega}$

Expresión explícita de la velocidad angular

$\underline{\omega}, \underline{\Omega}$ función de \underline{n}, ϕ y de $\underline{\dot{n}}, \dot{\phi}$

$$\mathbf{R}^T \left(\mathbf{R} - \mathbf{I} \right) \dot{\underline{n}} = \mathbf{L}^T \mathbf{R}^T \dot{\underline{n}} \quad \mathbf{R}^T \underline{n} = \underline{n} \quad \left(\mathbf{R}^T - \mathbf{I} \right) \dot{\underline{n}} = -\mathbf{R}^T \dot{\underline{n}}$$

$$\left(\mathbf{R}^T - \mathbf{I} \right) \dot{\underline{n}} = \mathbf{R}^T \dot{\underline{n}} = \underline{\Omega} \underline{n} \quad \left(\mathbf{R} - \mathbf{I} \right) \dot{\underline{n}} = \underline{\omega} \underline{n}$$

$$\underline{\dot{n}} \underline{\Omega} = \left(\mathbf{I} - \mathbf{R}^T \right) \dot{\underline{n}}$$

$$\underline{\dot{n}} \underline{\omega} = \left(\mathbf{R} - \mathbf{I} \right) \dot{\underline{n}}$$

Ver en el libro $\Rightarrow \underline{n}^T \underline{\Omega} = \dot{\phi}$

$\underline{n}^T \underline{\omega} = \dot{\phi}$

Luego:

$$\begin{bmatrix} 2c_{11} \\ s_{11} \end{bmatrix} \underline{\omega} = \begin{pmatrix} (I-R) \underline{s}_1 \\ \dot{\phi} \end{pmatrix}$$

$$\begin{bmatrix} 2c_{11} \\ s_{11} \end{bmatrix} \underline{\omega} = \begin{pmatrix} (R-I) \underline{n} \\ \dot{\phi} \end{pmatrix}$$

Resolviendo:

$$\underline{\omega} = \underbrace{\begin{bmatrix} 2c_{11} \\ s_{11} \end{bmatrix}}_{\underline{M}^T} (R-I) \underline{s}_1 + \underline{n} \dot{\phi} = \underline{M} \underline{s}_1 + \underline{n} \dot{\phi}$$

$$\underline{\Omega} = \underbrace{\begin{bmatrix} \underline{n}^2 \\ (I-R) \end{bmatrix}}_{\underline{M}} \underline{s}_1 + \underline{s}_1 \dot{\phi} = \underline{M}^T \underline{s}_1 + \underline{s}_1 \dot{\phi}$$

$$\underline{M} = \sin \phi \underline{I} - (1 - \cos \phi) \underline{n} \underline{n}^T$$

$$M = -\underline{\underline{n}} (\underline{\underline{I}} - \underline{\underline{P}}^T)$$

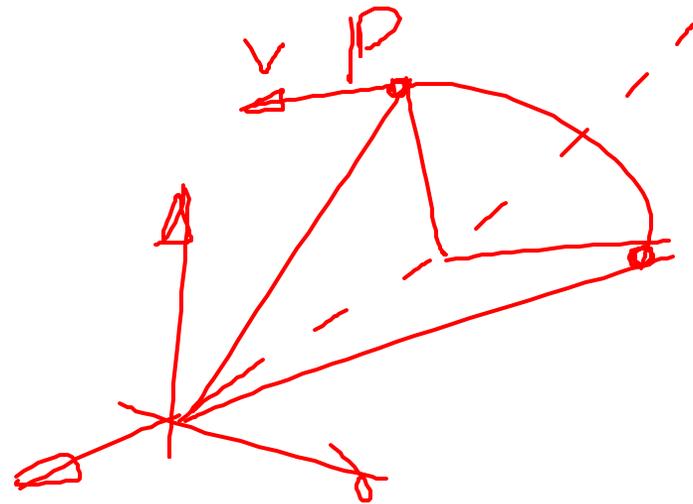
$$M^T = -(\underline{\underline{R}} - \underline{\underline{I}}) (\underline{\underline{n}})^T = +(\underline{\underline{R}} - \underline{\underline{I}}) \underline{\underline{n}} = \underline{\underline{n}} (\underline{\underline{R}} - \underline{\underline{I}})$$

$$\underline{\underline{R}} \underline{\underline{n}} = \underline{\underline{R}} \underline{\underline{n}} \quad \underline{\underline{P}}^T \underline{\underline{R}} = \underline{\underline{R}} \underline{\underline{n}} \quad \underline{\underline{R}} = \underline{\underline{n}} \underline{\underline{R}}$$

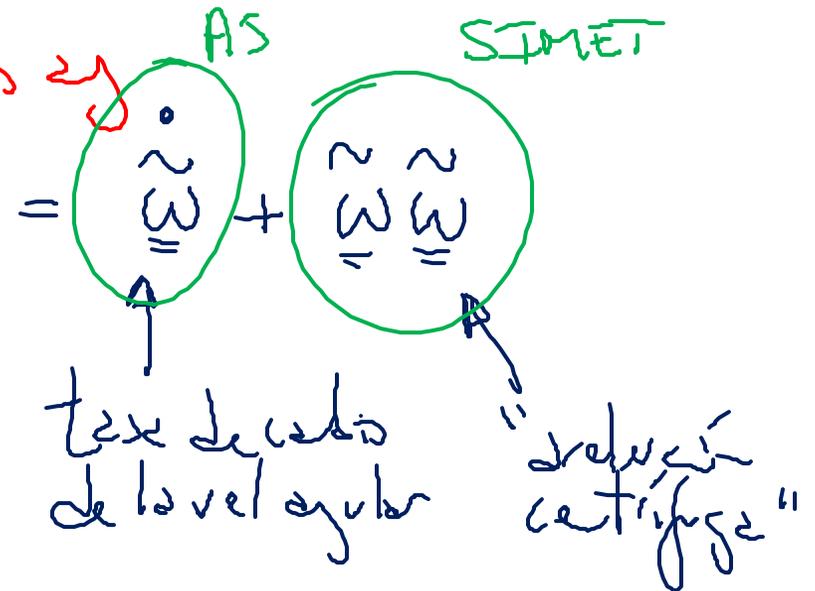
Aceleración zylar

$$a_p = \dot{v}_p = \ddot{x}_p = \ddot{R}X =$$

$$= \underbrace{\ddot{R}R^T}_{\text{matriz de aceleración zylar}} x$$



$$a = \ddot{R}R^T = \frac{d}{dt} \underbrace{\underbrace{(R R^T)}_{\tilde{\omega}}}_{\text{matriz de aceleración zylar}} - \underbrace{\underbrace{\dot{R}R^T}_{\tilde{\omega}}}_{\tilde{\omega}} \underbrace{\underbrace{R R^T}_{\tilde{\omega}}}_{\tilde{\omega}}$$



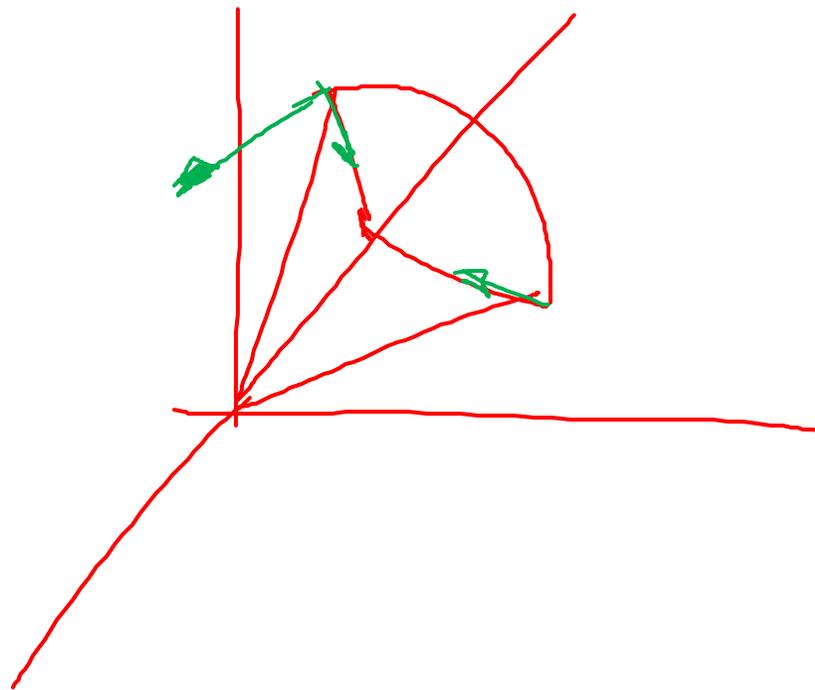
$$a_p = \underbrace{\tilde{\omega}}_{\text{rate of change of angular velocity}} x + \underbrace{\tilde{\omega}\tilde{\omega}}_{\text{centrifugal acceleration}} x$$

En coord materiales:

vel
del pto $\underline{A}_p = \underline{R}^T \underline{\dot{p}} = \underline{R}^T \underline{\dot{R}} X$

vel
angular $\underline{A} = \underline{R}^T \underline{\dot{R}} = \frac{d}{dt} (\underline{R}^T \underline{R}) - \underline{R}^T \underline{\dot{R}} = \underline{\Omega} + \underline{\Omega} \underline{\Omega}$

$$\underline{A} = \underline{R}^T \underline{\dot{R}}$$



Movimiento de rotación infinitesimal

$$\underline{x} = \underline{R} \underline{X}$$

$$\dot{\underline{x}} = \dot{\underline{R}} \underline{X}$$

$$\delta \underline{x} = \delta \underline{R} \underline{X} \quad \text{desplaza virtuales}$$

$$= \delta \underline{R} \underline{R}^T \underline{x} = \delta \underline{\theta} \underline{x}$$

$$\delta \underline{x} = \underline{R} \underline{R}^T \delta \underline{R} \underline{X} = \underline{R} \delta \underline{\theta} \underline{X}$$

$$\delta \underline{\theta} = \delta \underline{R} \underline{R}^T$$

$$\delta \underline{\theta} = \underline{R}^T \delta \underline{R}$$

\Downarrow

$$\delta \underline{R} = \delta \underline{\theta} \underline{R} = \underline{R} \delta \underline{\theta}$$

$$\underline{R} \delta \underline{\theta} = \delta \underline{\theta}$$

Variación de velocidades angulares

$$\underline{\tilde{\Omega}} = \underline{R}^T \underline{\dot{R}} \Rightarrow \delta \underline{\tilde{\Omega}} = \delta \underline{R}^T \underline{\dot{R}} + \underline{R}^T \delta \underline{\dot{R}}$$

$$\delta \underline{\Theta} = \underline{R}^T \delta \underline{R} \Rightarrow \frac{d}{dt} (\delta \underline{\Theta}) = \underline{\dot{R}}^T \delta \underline{R} + \underline{R}^T \delta \underline{\dot{R}}$$

Luego:

$$\delta \underline{\tilde{\Omega}} = \delta \underline{\dot{\Theta}} + \delta \underline{R}^T \underline{\dot{R}} - \underline{\dot{R}}^T \delta \underline{R} = \delta \underline{\dot{\Theta}} - \delta \underline{\Theta} \underline{\tilde{\Omega}} + \underline{\tilde{\Omega}} \delta \underline{\Theta}$$

$$\delta \underline{\tilde{\Omega}} = \delta \underline{\dot{\Theta}} + \delta \underline{\Theta} \underline{\tilde{\Omega}} - \underline{\tilde{\Omega}} \delta \underline{\Theta}$$

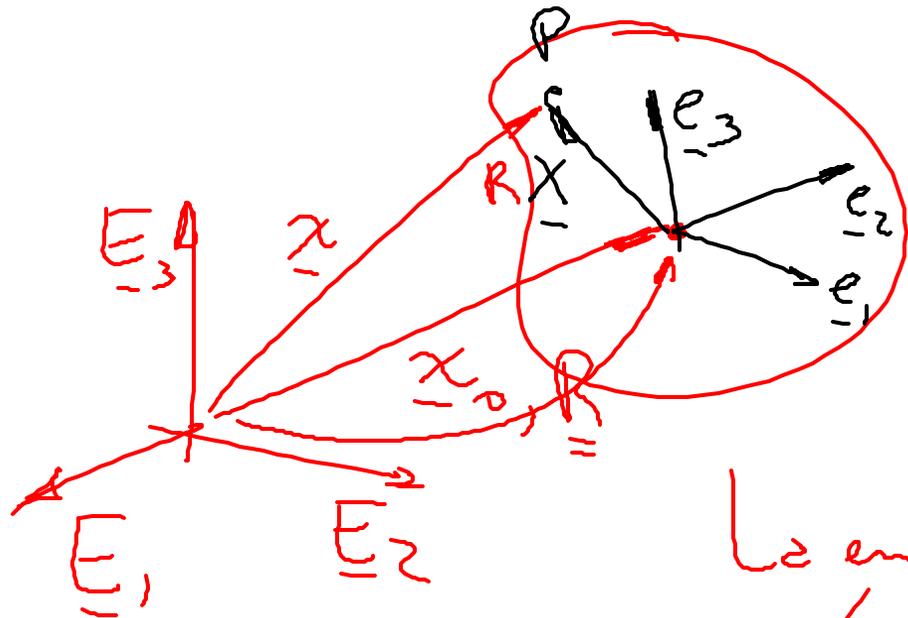
$$\delta \underline{\tilde{\Omega}} = \delta \underline{\dot{\Theta}} - \underbrace{\delta \underline{\tilde{\Theta}} \underline{\tilde{\Omega}} + \underline{\tilde{\Omega}} \delta \underline{\tilde{\Theta}}}$$

$$\begin{aligned} \left[\underline{\tilde{a}} \underline{\tilde{b}} - \underline{\tilde{b}} \underline{\tilde{a}} \right] c &= \underline{a} \times (\underline{b} \times c) - \underline{b} \times (\underline{a} \times c) = \\ &= (\underline{a} \cdot c) \underline{b} - \cancel{(\underline{a} \cdot \underline{b}) c} - (\underline{b} \cdot c) \underline{a} + \cancel{(\underline{b} \cdot \underline{a}) c} = \\ &= c \times (\underline{b} \times \underline{a}) = (\underline{a} \times \underline{b}) \times c = \underline{(\tilde{a} \times \tilde{b})} c \end{aligned}$$

$$\left[\underline{\tilde{a}} \underline{\tilde{b}} - \underline{\tilde{b}} \underline{\tilde{a}} \right] = \underline{\tilde{a} \times \tilde{b}}$$

$$\delta \underline{\tilde{\Omega}} = \delta \underline{\dot{\Theta}} + \underline{(\tilde{\Omega} \times \delta \Theta)} \Rightarrow \frac{\delta \underline{\Omega} = \delta \underline{\dot{\Theta}} + \underline{\tilde{\Omega}} \delta \underline{\Theta}}{\delta \underline{\omega} = \delta \underline{\dot{\Theta}} - \underline{\tilde{\omega}} \delta \underline{\Theta}}$$

CUERPO RIGIDO



$$\underline{e}_i = \underline{R}(t) \underline{E}_i$$

(en el instante inicial, $\underline{e}_i \equiv \underline{E}_i$)

$$\underline{x} = \underline{x}_0 + \underline{R}(t) \underline{X}$$

La energía cinética:

$$K = \frac{1}{2} \left(m \underline{\dot{x}}_0^T \underline{\dot{x}}_0 + \underline{\Omega}^T \underline{J} \underline{\Omega} \right)$$

de porque usas
vel ang matrix

$$K = \frac{1}{2} \int_{\mathcal{V}} \dot{\underline{x}}^T \dot{\underline{x}} \rho \, dV = \frac{1}{2} \int_{\mathcal{V}} \dot{\underline{x}}_0^T \dot{\underline{x}}_0 \rho \, dV + \frac{1}{2} \underbrace{\int_{\mathcal{V}} (\dot{R}X)^T (\dot{R}X) \rho \, dV}_{(*)}$$

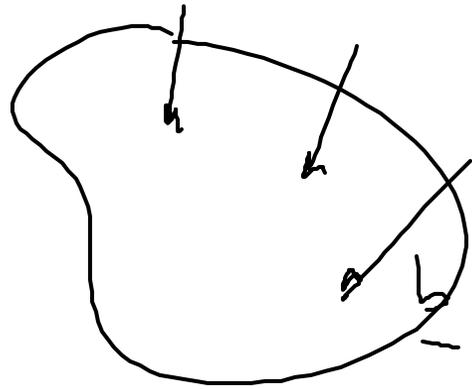
$$\dot{R}X = R \tilde{\Omega} X = -R \tilde{X} \Omega \quad \Rightarrow \quad \tilde{\omega} R X = -\tilde{R} X \omega = -R \tilde{X} R^T \omega$$

$$(*) = \frac{1}{2} \int_{\mathcal{V}} \Omega^T X R^T R \tilde{X} \Omega \rho \, dV = \frac{1}{2} \Omega^T \underbrace{\int_{\mathcal{V}} \tilde{X}^T X \rho \, dV}_{\underline{Y}} \Omega$$

$$K = \frac{1}{2} \left(m \dot{\underline{x}}_0^T \dot{\underline{x}}_0 + \underline{\Omega}^T \underline{Y} \underline{\Omega} \right) \quad \underline{I}(t) = \underline{R} \underline{Y} \underline{R}^T$$

$$K = \frac{1}{2} \left(m \dot{\underline{x}}_0^T \dot{\underline{x}}_0 + \underline{\omega}^T \underline{I}(t) \underline{\omega} \right)$$

CARGAS EXTERNAS DEVOLUTIVAS



$$P = + \int \underline{x}^T \underline{b} dV =$$

$$+ \int (\underline{x}_0 + R\underline{x})^T \underline{b} dV$$

$$\therefore P = + \underline{x}_0^T \underline{f} + \int (R\underline{x})^T \underline{b} dV$$

$$\delta P = + \delta \underline{x}_0^T \underline{f} + \delta \underline{\Theta}^T \underline{C} \quad \underline{C} = \int \underline{\tilde{x}}^T \underline{R}^T \underline{b} dV$$

ECS DE MOVIMIENTO

$$\delta \int_{t_1}^{t_2} \mathcal{L} dt = 0 = \delta \int_{t_1}^{t_2} (\mathcal{K} - \mathcal{P}) dt$$

$$\delta \mathcal{K} = \delta \underline{\dot{x}}_0^T m \underline{\dot{x}}_0 + \delta \underline{\Omega}^T \underline{J} \underline{\Omega} =$$

$$= \delta \underline{\dot{x}}_0^T m \underline{\dot{x}}_0 + \left(\delta \underline{\dot{\theta}} + \underline{\tilde{\Omega}} \delta \underline{\theta} \right)^T \underline{J} \underline{\Omega}$$

$$\int_{t_1}^{t_2} \delta \mathcal{K} dt = \cancel{\delta \underline{x}_0^T m \underline{x}_0 \Big|_{t_1}^{t_2}} - \int_{t_1}^{t_2} \delta \underline{x}_0^T m \underline{\ddot{x}}_0 dt + \cancel{\delta \underline{\theta}^T \underline{J} \underline{\Omega} \Big|_{t_1}^{t_2}} - \int_{t_1}^{t_2} \delta \underline{\theta}^T \underline{J} \underline{\dot{\Omega}} dt - \int_{t_1}^{t_2} \delta \underline{\theta}^T \underline{\tilde{\Omega}} \underline{J} \underline{\Omega} dt$$

$$\delta \int_{t_1}^{t_2} K dt = \int_{t_1}^{t_2} \delta \underline{x}_0^T m \dot{\underline{x}}_0 dt + \int_{t_1}^{t_2} \delta \underline{\theta}^T \left(\underline{J} \dot{\underline{\Omega}} + \underline{\tilde{\Omega}} \underline{J} \underline{\Omega} \right) dt$$

$$\delta \int_{t_1}^{t_2} (-\mathcal{P}) dt = \int_{t_1}^{t_2} \delta \underline{x}_0^T \underline{f} + \delta \underline{\theta}^T \underline{C}$$

...

$$m \dot{\underline{x}}_0 = \underline{f}$$

$$\underline{J} \dot{\underline{\Omega}} + \underline{\tilde{\Omega}} \underline{J} \underline{\Omega} = \underline{C}$$

Pág 114, Landau \Rightarrow Ecuaciones de Euler

La relación más simple entre aspectos del movimiento angular rotacional y los comp de la vel angular ocurre en un sist de coord móvil con el cuerpo, en el que \underline{P} de tal los ejes coincide con los ejes ppals de inercia.

Sea $\frac{d\underline{A}}{dt}$ la tasa de cambio de un vector \underline{A} respecto al sist inercial. Si \underline{A} no cambia en un sistema móvil, su tasa de cambio es sólo debida a la rotación y entonces

$$\frac{d\underline{A}}{dt} = \underline{\Omega} \times \underline{A}$$

En un caso gen, tenes también el cambio de \underline{A} y resulta:

$$\frac{d\underline{A}}{dt} = \frac{d'\underline{A}}{dt} + \underline{\Omega} \times \underline{A}$$

cambio de \underline{A} en el sist. móvil

Usando estas formulas P/ el momento lineal y el angular:

$$\left(\frac{d'\underline{P}}{dt} + \underline{\Omega} \times \underline{P} = \underline{F} \right)$$

$$\underline{P} = \underline{\mu} \underline{V}$$

$$\left(\frac{d'\underline{M}}{dt} + \underline{\Omega} \times \underline{M} = \underline{C} \right)$$

$$\underline{M} = \underline{J} \underline{\Omega}$$

$$\begin{aligned} M_1 &= J_1 \Omega_1 \\ M_2 &= J_2 \Omega_2 \\ M_3 &= J_3 \Omega_3 \end{aligned}$$

$$\mu \left(\frac{dV_1}{dt} + \Omega_2 V_3 - \Omega_3 V_2 \right) = F_1$$

$$\mu \left(\frac{dV_2}{dt} + \Omega_3 V_1 - \Omega_1 V_3 \right) = F_2$$

$$\mu \left(\frac{dV_3}{dt} + \Omega_1 V_2 - \Omega_2 V_1 \right) = F_3$$

$$J_1 \frac{d\Omega_1}{dt} + (J_3 - J_2) \Omega_2 \Omega_3 = C_1$$

$$J_2 \frac{d\Omega_2}{dt} + (J_1 - J_3) \Omega_1 \Omega_3 = C_2$$

$$J_3 \frac{d\Omega_3}{dt} + (J_2 - J_1) \Omega_1 \Omega_2 = C_3$$

Tiempo simétrico:

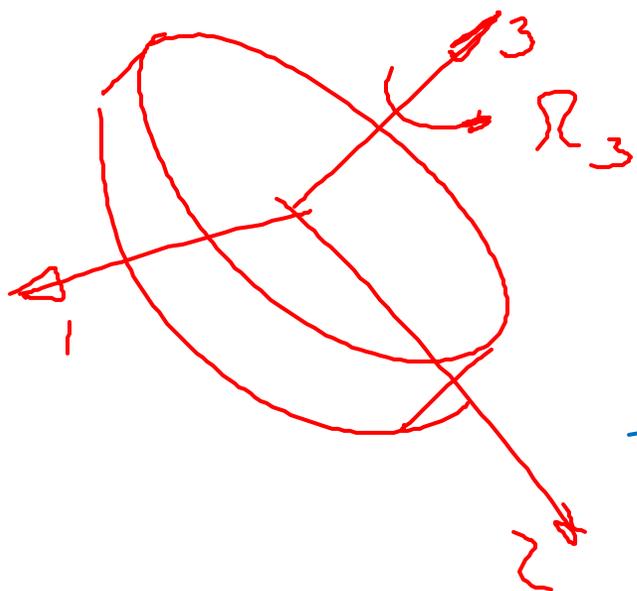
$$J_1 = J_2$$

$$J_1 = J_2$$

$$C_1 = C_2 = C_3 = 0$$

$$\Rightarrow \frac{d\Omega_3}{dt} = 0$$

$$\Omega_3 = \text{cte}$$



$$\frac{d\Omega_1}{dt} + \frac{J_3 - J_2}{J_1} \Omega_2 \Omega_3 = 0$$

$$\frac{d\Omega_1}{dt} + \omega \Omega_2 = 0$$

$$+ i \left(\frac{d\Omega_2}{dt} - \omega \Omega_1 \right) = 0$$

$$\frac{d}{dt} (\Omega_1 + i\Omega_2) = i\omega (\Omega_1 + i\Omega_2)$$

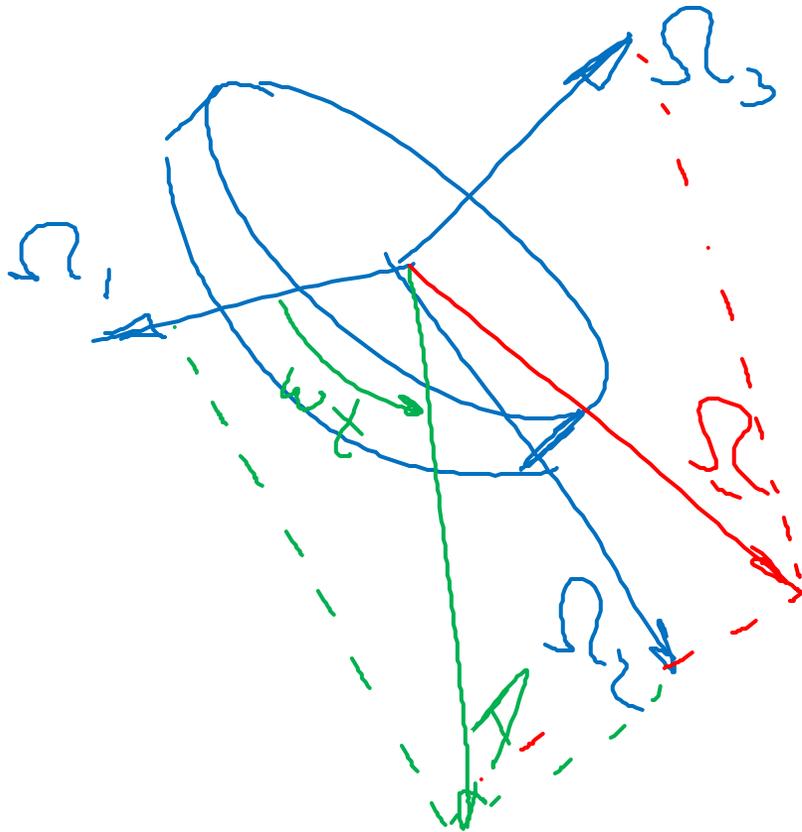
$$\Omega_1 + i\Omega_2 = A \exp(i\omega t)$$

$$A = \text{cte}$$

$$\left\{ \begin{array}{l} \Omega_1 = A \cos \omega t \\ \Omega_2 = A \sin \omega t \\ \Omega_3 = cte \end{array} \right\}$$

$$\omega = \Omega_3 \frac{J_3 - J_1}{J_1}$$

$$A = \sqrt{\Omega_1^2 + \Omega_2^2}$$

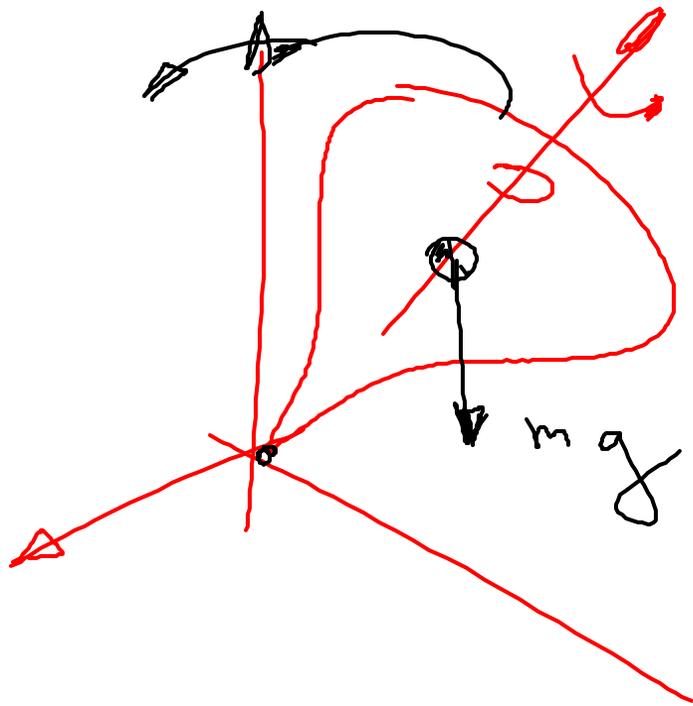


CI?

$$\left. \begin{array}{l} \text{i) } \Omega_3 = cte \\ \Omega_1 = \Omega_2 = 0 \end{array} \right\}$$

$$\left. \begin{array}{l} \Omega_3 = cte \\ \Omega_1 = \Omega_2 = 0 \end{array} \right\}$$

$$\left. \begin{array}{l} \text{ii) } \Omega_1 \\ \Omega_2 = 0 \end{array} \right\}$$



PARAMETRIZACION DE ROTACIONES

$$\underline{R} = \underline{R}(\underline{a}) \quad \underline{a} = (\alpha_1, \alpha_2, \alpha_3)$$

$$\underline{R} = \underline{R}(n, \phi)$$

$$\underline{\omega} = \underline{P}(\underline{a}) \dot{\underline{a}} \quad \underline{\Omega} = \underline{Q}(\underline{a}) \dot{\underline{a}}$$

$$\underline{\omega} = \underline{R} \underline{\Omega} \quad \Rightarrow \quad \underline{R} = \underline{P} \underline{Q}^{-1}$$

i) Vector de rotación Cartesiano (CRV)

$$\underline{\Psi} = \underline{n} \phi$$

coord mínimo (3)
libre de singularidades

$$\underline{R} = \underline{I} + \frac{\sin \|\underline{\Psi}\|}{\|\underline{\Psi}\|} \underline{\Psi} + \frac{1 - \cos \|\underline{\Psi}\|}{\|\underline{\Psi}\|^2} \underline{\Psi} \underline{\Psi}^T$$

$$\underline{R} = \exp(\underline{\hat{\Psi}})$$

$$\left(\underline{R} = \underline{I} \underline{I}^T \right)$$

$$\underline{\Omega} = \underline{T}(\underline{\Psi}) \dot{\underline{\Psi}} \quad \underline{\omega} = \underline{T}^T(\underline{\Psi}) \dot{\underline{\Psi}}$$

$$\underline{T}(\underline{\Psi}) = \underline{I} + \left(\frac{\cos \|\underline{\Psi}\| - 1}{\|\underline{\Psi}\|^2} \right) \underline{\Psi} \underline{\Psi}^T + \left(\frac{\sin \|\underline{\Psi}\|}{\|\underline{\Psi}\|} \right) \underline{\hat{\Psi}}$$

$$\lim_{\Psi \rightarrow 0} T(\Psi) = \lim_{\Psi \rightarrow 0} R(\Psi) = \underline{\underline{I}}$$

rot "moderadas"

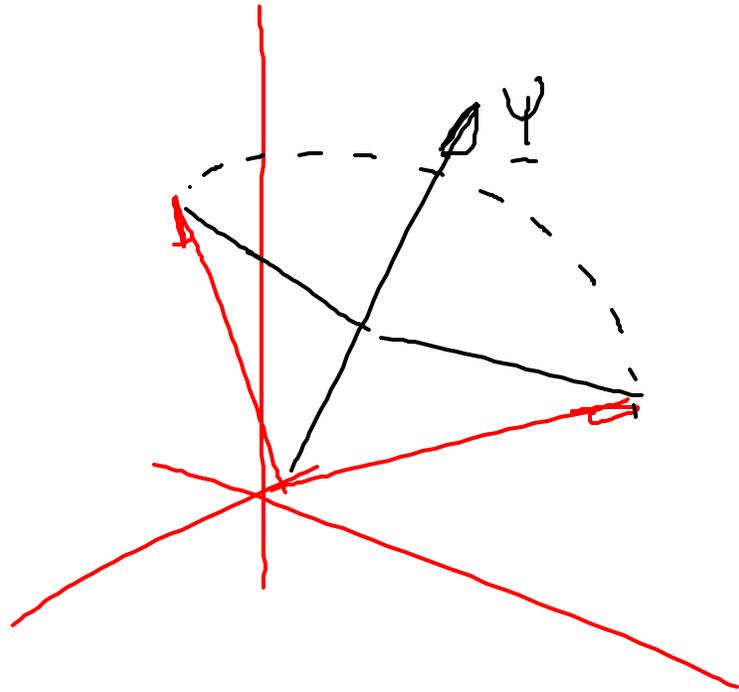
$$R(\Psi) = \exp(\tilde{\Psi}) \approx \left(I + \tilde{\Psi} + \frac{1}{2} \tilde{\Psi}^2 \right) + \frac{1}{3!} \tilde{\Psi}^3 \dots$$

$$T(\Psi) = I - \frac{1}{2!} \tilde{\Psi} + \frac{1}{3!} \tilde{\Psi}^2 \dots$$

Aceleraciones angulares:

$$\begin{aligned} \overset{\circ}{\Omega} &= T(\Psi) \overset{\circ}{\dot{\Psi}} + \dot{T}(\Psi) \overset{\circ}{\Psi} \\ \overset{\circ}{\omega} &= T^T(\Psi) \overset{\circ}{\dot{\Psi}} + \dot{T}^T(\Psi) \overset{\circ}{\Psi} \end{aligned}$$

$$\begin{aligned} \delta \Theta &= T(\Psi) \delta \Psi \\ \delta \theta &= T^T(\Psi) \delta \Psi \end{aligned}$$



Parámetros de Rodrigues

$$\underline{b} = n \tan \frac{\phi}{2}$$

$$\underline{R} = \underline{I} + \frac{2}{1 + \|\underline{b}\|^2} \left(\underline{\underline{b}} + \underline{\underline{b}} \underline{\underline{b}} \right)$$

$$\underline{\Omega} = \underline{T}(\underline{b}) \dot{\underline{b}} \quad \underline{\omega} = \underline{T}^T(\underline{b}) \dot{\underline{b}}$$

$$\underline{T} = \frac{2}{1 + \|\underline{b}\|^2} \left(\underline{I} - \underline{\underline{b}} \right)$$

par - vimo (3)
oper - lineales
Singularidad en $\pm \pi$

Parámetros de Euler

INCONV # puntos > 3

vent 4 - dimensional

$$\underline{P}^T = [\underline{e}_0 \underline{e}^T] = [e_0 e_1 e_2 e_3]$$

$$e_0 = \cos \phi/2$$

$$\underline{e} = \underline{n} \sin \phi/2$$

$$\Rightarrow e_0^2 + e_1^2 + e_2^2 + e_3^2 = 1$$

$$\underline{R} = (2e_0^2 - 1) \underline{I} + 2\underline{e}\underline{e}^T + 2e_0 \underline{\tilde{e}}$$

$$\underline{R} = \underline{H} \underline{G}^T$$

$$\underline{H} = \begin{bmatrix} -\underline{e} & e_0 \underline{I} + \underline{\tilde{e}} \end{bmatrix}$$
$$\underline{G} = \begin{bmatrix} -\underline{e} & e_0 \underline{I} - \underline{\tilde{e}} \end{bmatrix}$$

$$\underline{\Omega} = 2\underline{G} \dot{\underline{p}}^0$$
$$\underline{\omega} = 2\underline{H} \dot{\underline{p}}^0$$

Algebra de quaternions

Def: quaternion \rightarrow n \acute{u} -coplejo de 4 dimensiones

$$\hat{q} = q_0 + i q_1 + j q_2 + k q_3 = q_0 + \underline{q}$$

i, j, k : n \acute{u} m imaginarios

$$i^2 = j^2 = k^2 = -1$$

$$jk = -kj = i$$

$$ki = -ik = j$$

$$ij = -ji = k$$

Regla multiplicativa

$$\hat{r} = \hat{p} \hat{q} = \underbrace{p_0 q_0 - p_1 q_1}_{r_0} + \underbrace{p_0 q_1 + q_0 p_1 + p_2 q_3}_{r_1}$$

Nota: $\hat{p} \hat{q} \neq \hat{q} \hat{p}$

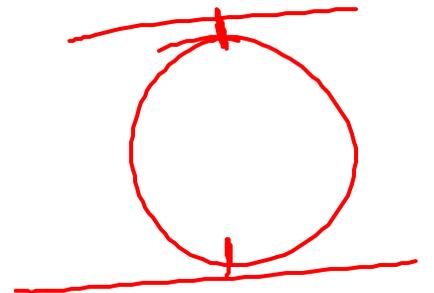
Quaternión conjugado

$$\hat{q}^* = q_0 - i q_1 - j q_2 - k q_3 = q_0 - \hat{q}$$

Norma:

$$\|\hat{q}\|^2 = \hat{q} \hat{q}^* = q_0^2 + q_1^2 + q_2^2 + q_3^2$$

		q_0
p_0	p_1	r_0
p_2	p_3	r_1



Rotación unitaria

$$\|\hat{e}\| = 1$$

Relación / Rotación finita

Sea $\hat{e} = \underline{e}_0 + \underline{e}$ quasi unitario

$\hat{X} = 0 + \underline{X}$ quasi vectorial

Se puede verificar que la rotación finita de \underline{X} a la posición \underline{x} correspondiente a parámetros de Euler $\hat{e} = \cos\alpha + \underline{n} \sin\alpha$

$$\hat{x} = \hat{e} \hat{X} \hat{e}^*$$

Nota: la rotación inversa

$$\hat{X} = \hat{e}^* \hat{x} \hat{e}$$

Comp de rotaciones:

$$\hat{x}_1 = \hat{e}_1 \hat{X} \hat{e}_1^*$$

$$\hat{x}_2 = \hat{e}_2 \hat{x}_1 \hat{e}_2^* = \underbrace{\hat{e}_2 \hat{e}_1}_{\hat{e}} \hat{X} \underbrace{\hat{e}_1^* \hat{e}_2^*}_{\hat{e}^*} = \hat{e} \hat{X} \hat{e}^*$$

$$\hat{e} = \hat{e}_2 \hat{e}_1$$

$$\hat{x} = \hat{e} \hat{X} \hat{e}^*$$

$$\hat{e}^* \hat{x} \hat{e} = \underbrace{\hat{e}^* \hat{e}}_1 \hat{X} \underbrace{\hat{e}^* \hat{e}}_1$$

$$\delta \underline{\Theta} \cdot \underline{f}(\underline{R}) = 0$$

$$\underline{R}(\underline{\Psi})$$

$$\delta \underline{\Psi} \cdot \underline{T}^T(\underline{\Psi}) \frac{\partial \underline{f}}{\partial \underline{\Theta}} \underline{T}(\underline{\Psi}) \Delta \underline{\Psi} = -R_{es}$$

$\underline{\Psi}_{inc}$

$$\underline{\Psi} = \underline{\Psi}_0 \circ \frac{\underline{\Psi}_{inc}}{inc\acute{o}g}$$

Otras opciones

Conformal Rotation Vector:

$$\underline{c} = 4 \underline{n} \tan \phi / 4$$

$$c_i = \frac{4e_i}{1+e_0}$$

∇ singularidad en $(-\pi, \pi)$

$$\underline{R} = \frac{1}{(4-c_0)^2} \left[(c_0^2 + 8c_0 - 16) \underline{I} + 2c_0 \underline{c} \underline{c}^T + 2c_0 \underline{\tilde{c}} \right]$$

Descripciones geométricas

Ang Euler

sewenz de 3 rot elementals:

a) rot (X_3, ψ)

b) rot (X_1, θ)

c) rot (X_3, ϕ)

\exists singularidad!

Ang Bryant

roll, pitch, yaw (navegar)

